

Figure 1.1: The derivative is the slope of the tangent line.

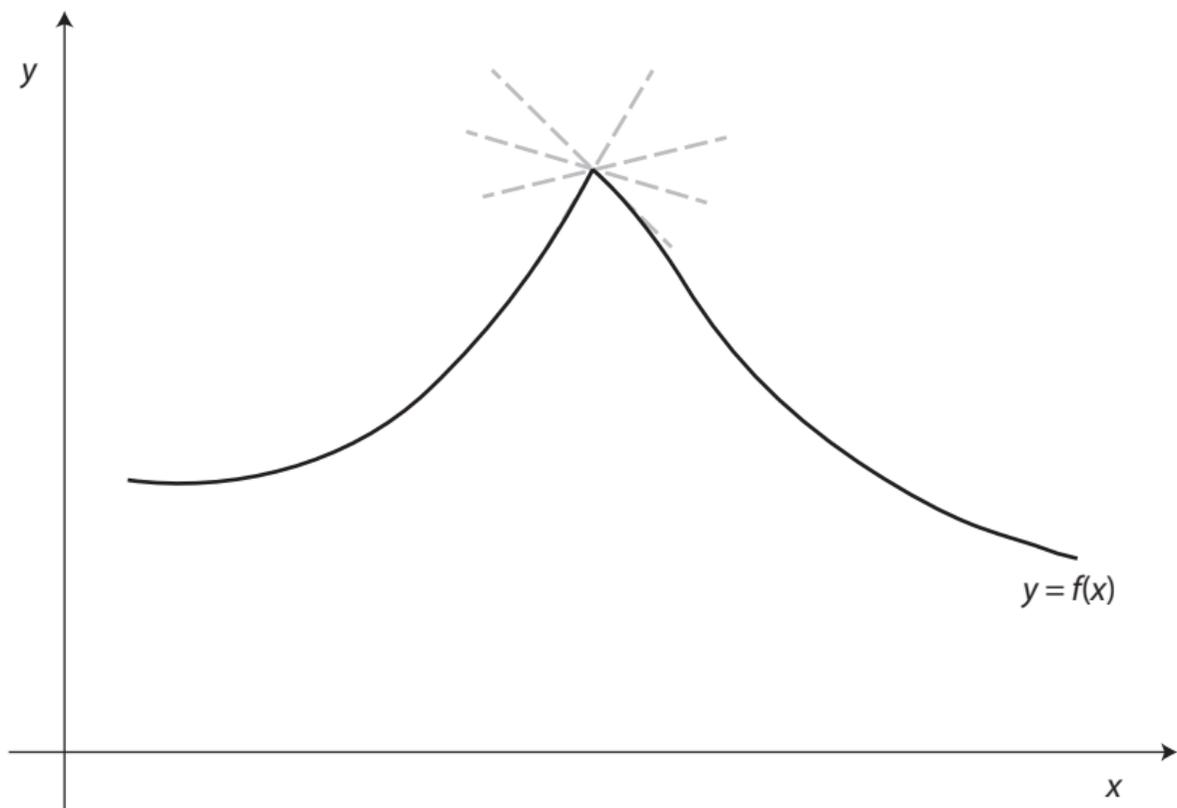


Figure 1.2: A graph with a point where there is no unique tangent line.

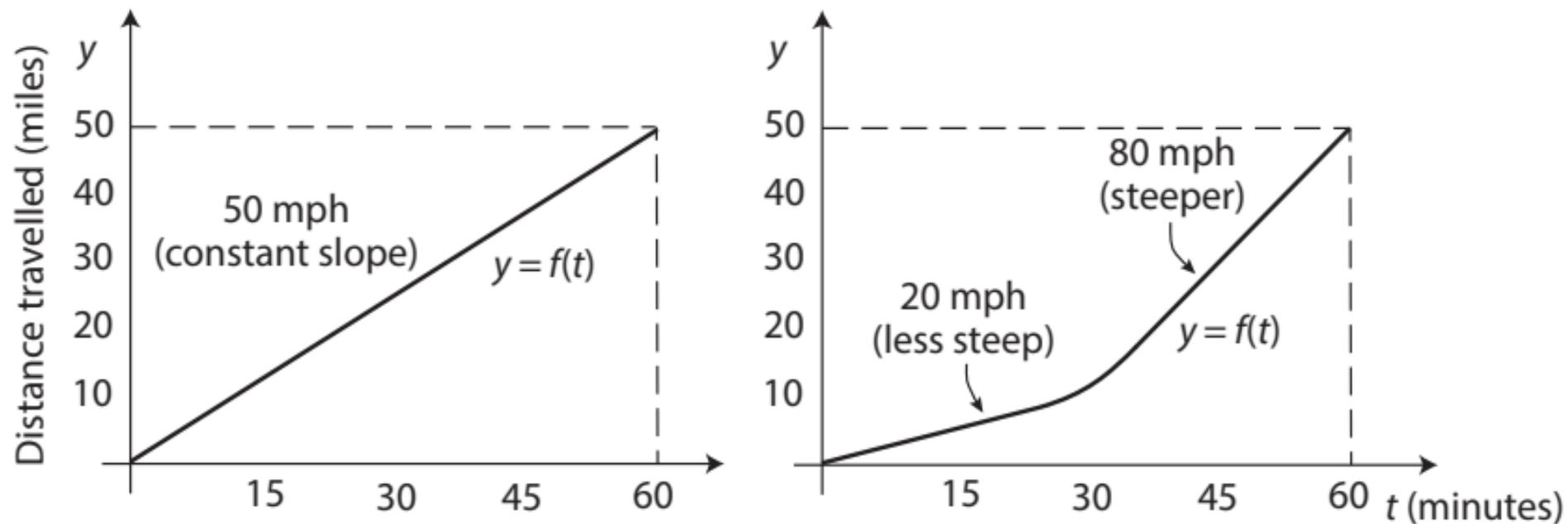


Figure 1.3: Two ways to travel 50 miles in an hour.

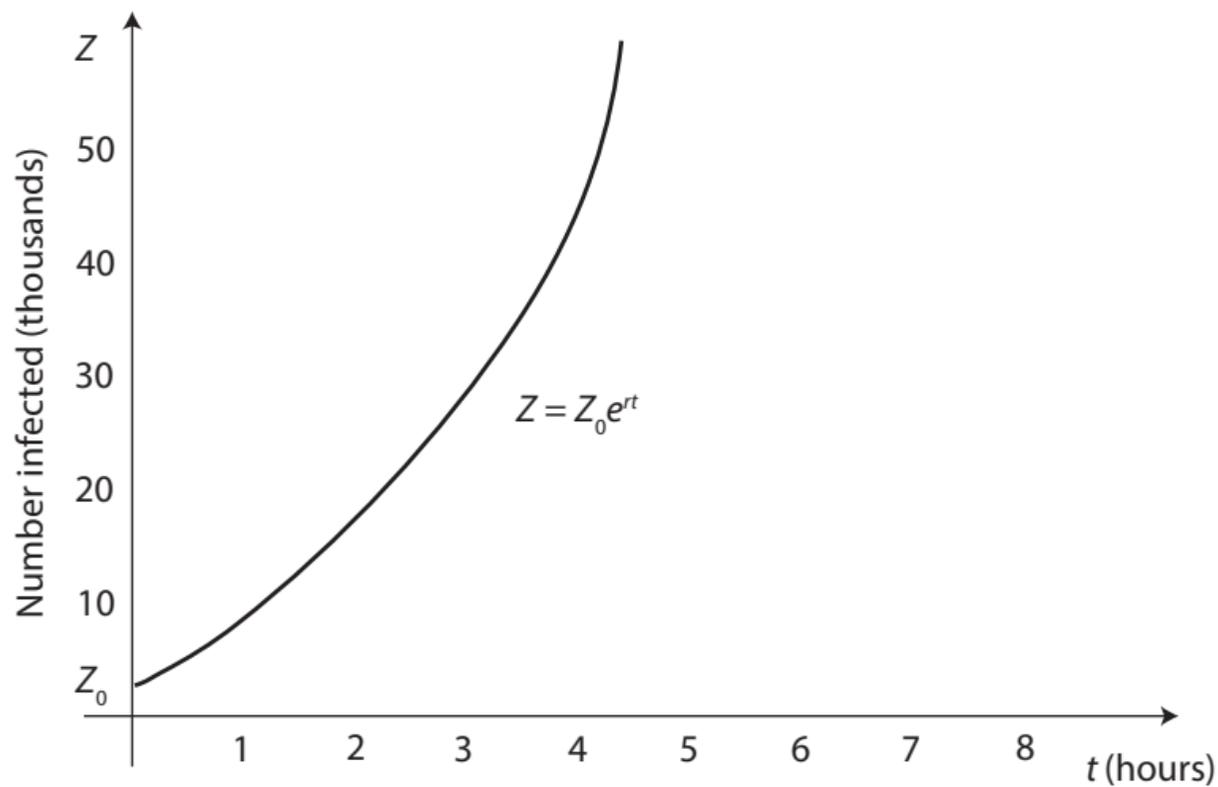


Figure 2.1: Exponential growth.

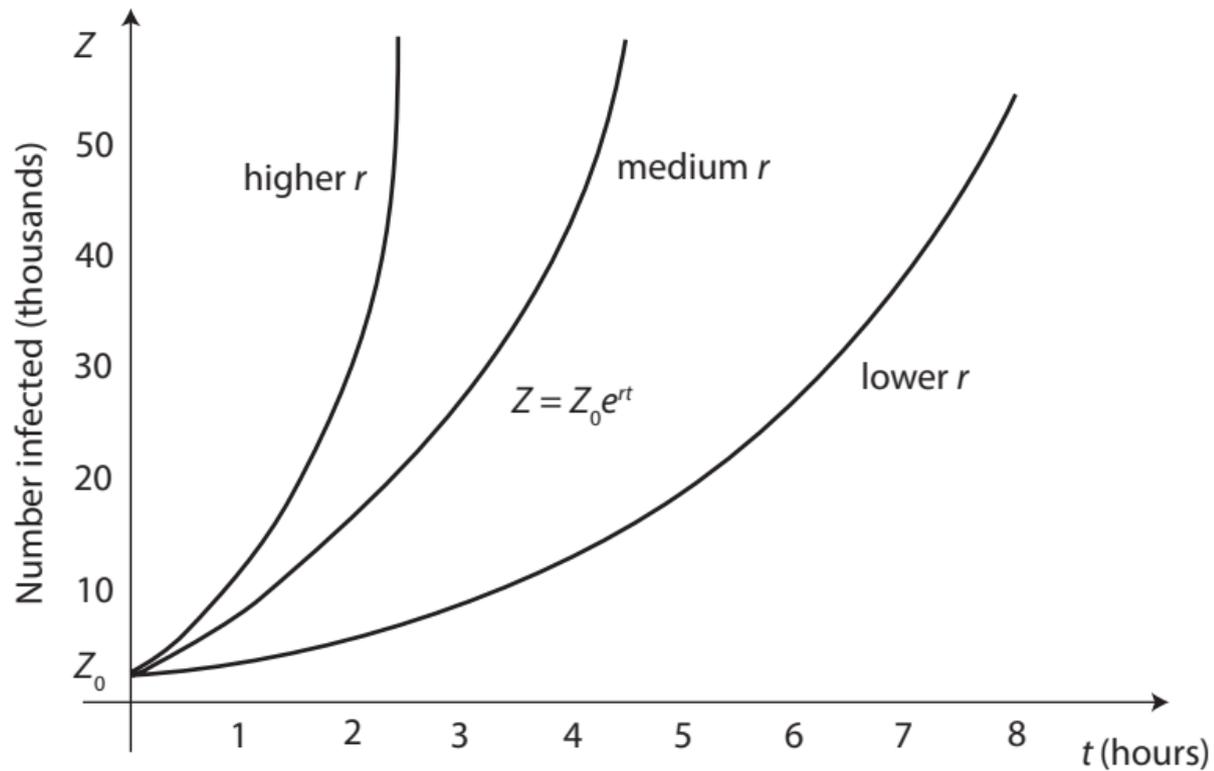


Figure 2.2: Growth rates yield different functions.

CHAPTER 3

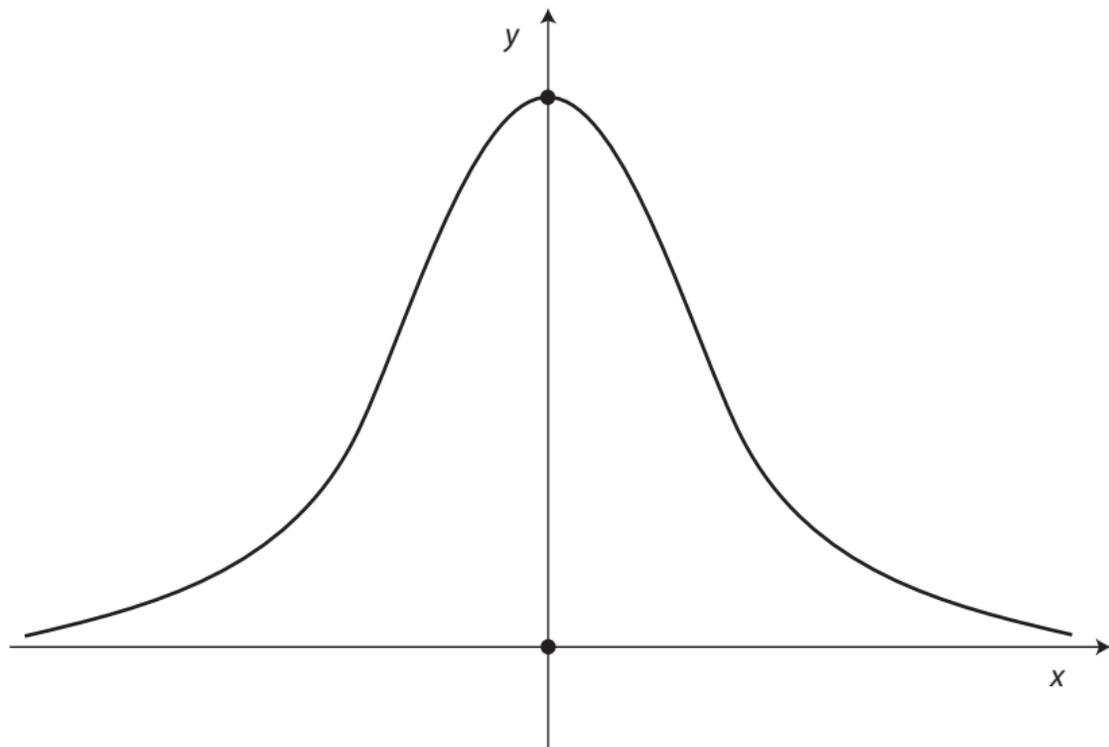


Figure 3.1: The standard normal distribution (bell curve).

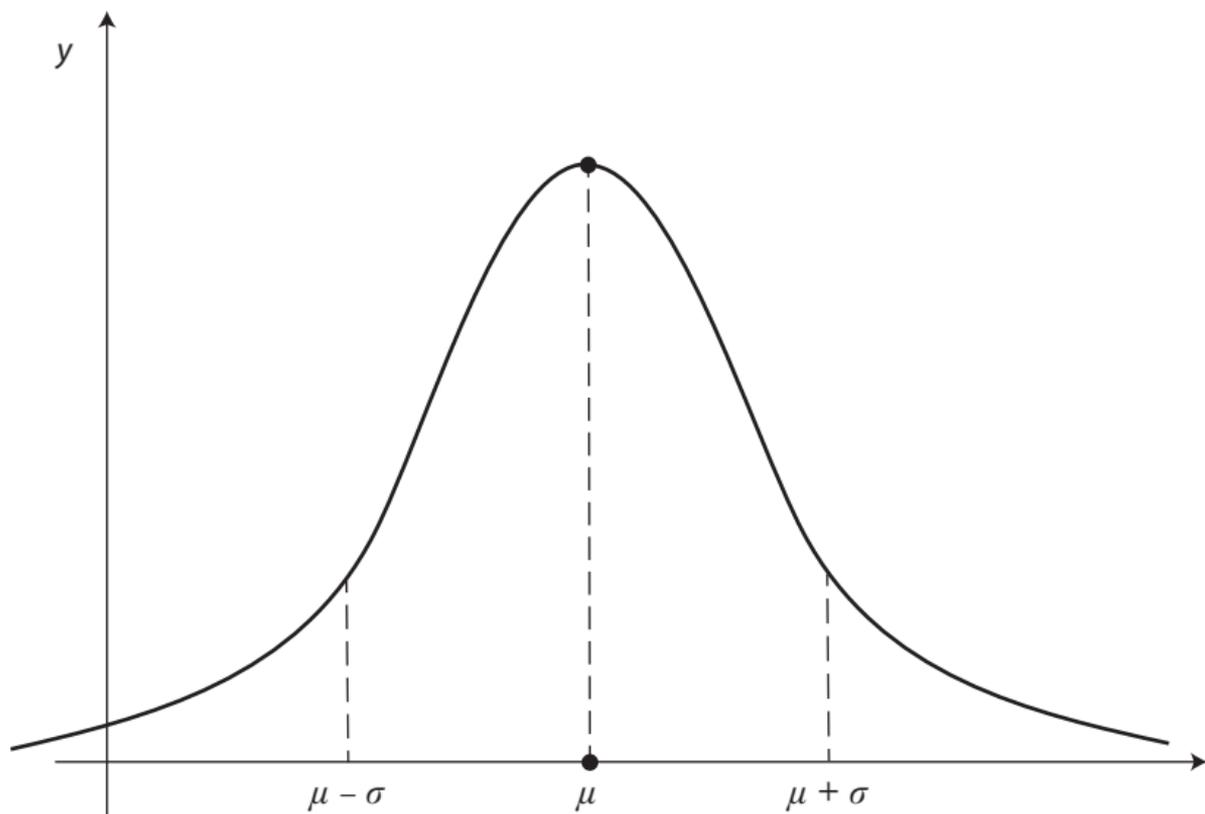


Figure 3.2: The general normal distribution.

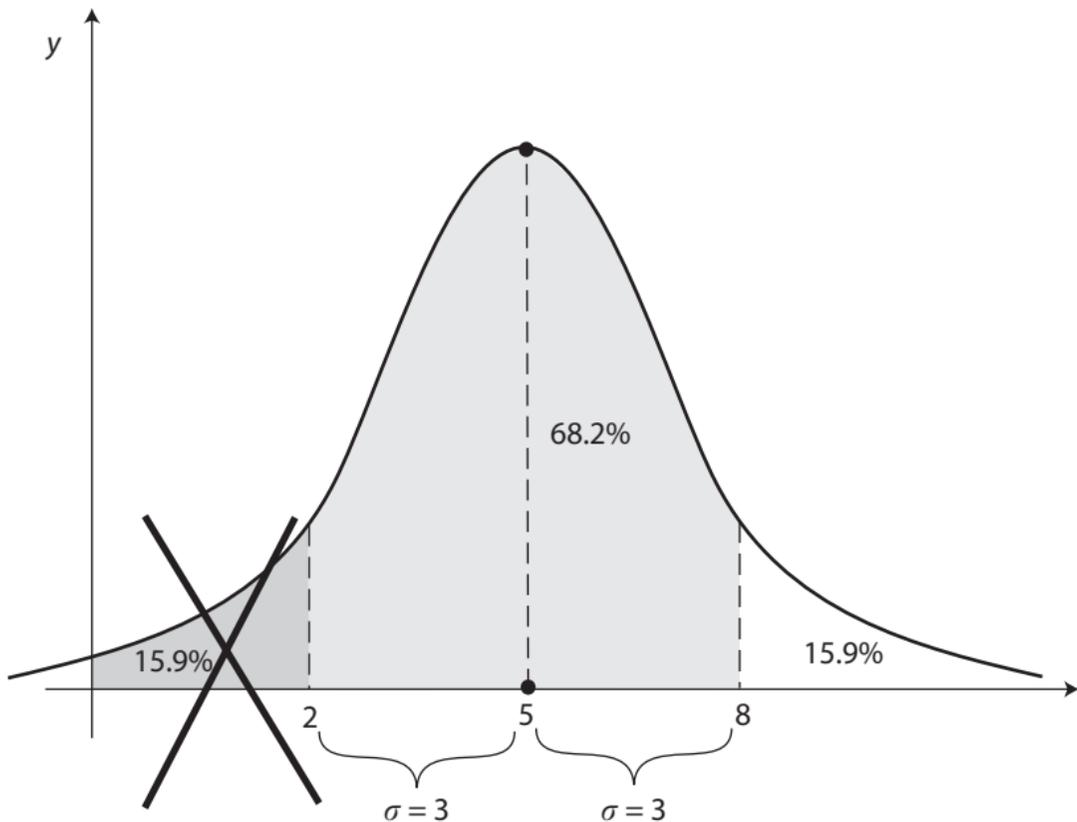


Figure 3.3: The people who run slower than the zombies.

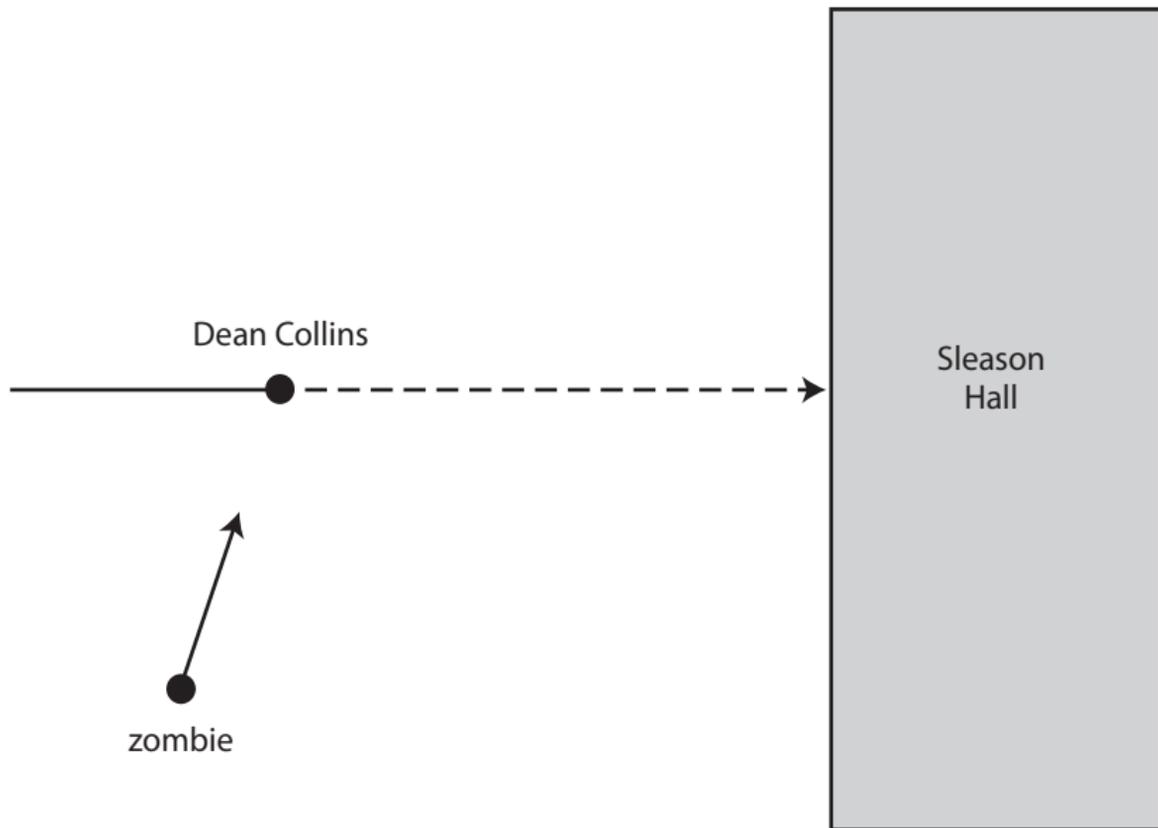


Figure 4.1: The dean's path and the zombie's path.

CHAPTER 4

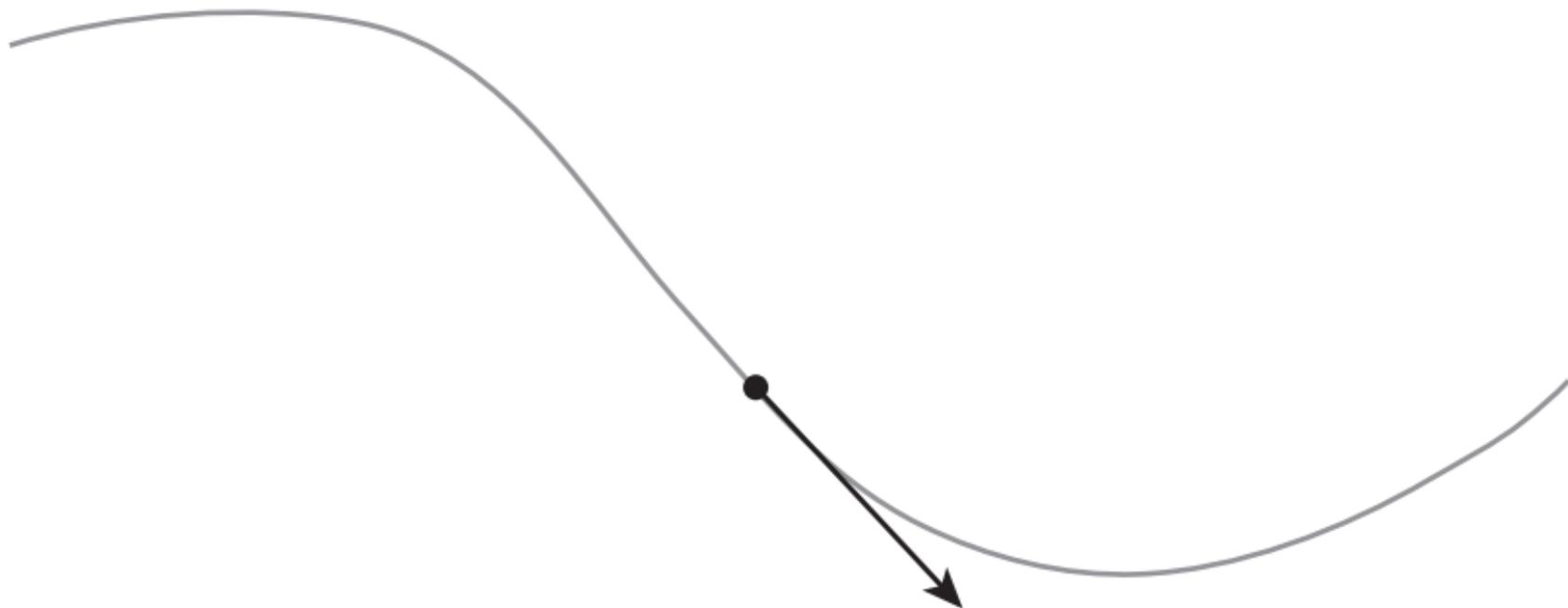


Figure 4.2: The tangent vector to a curve always points in the direction of motion.

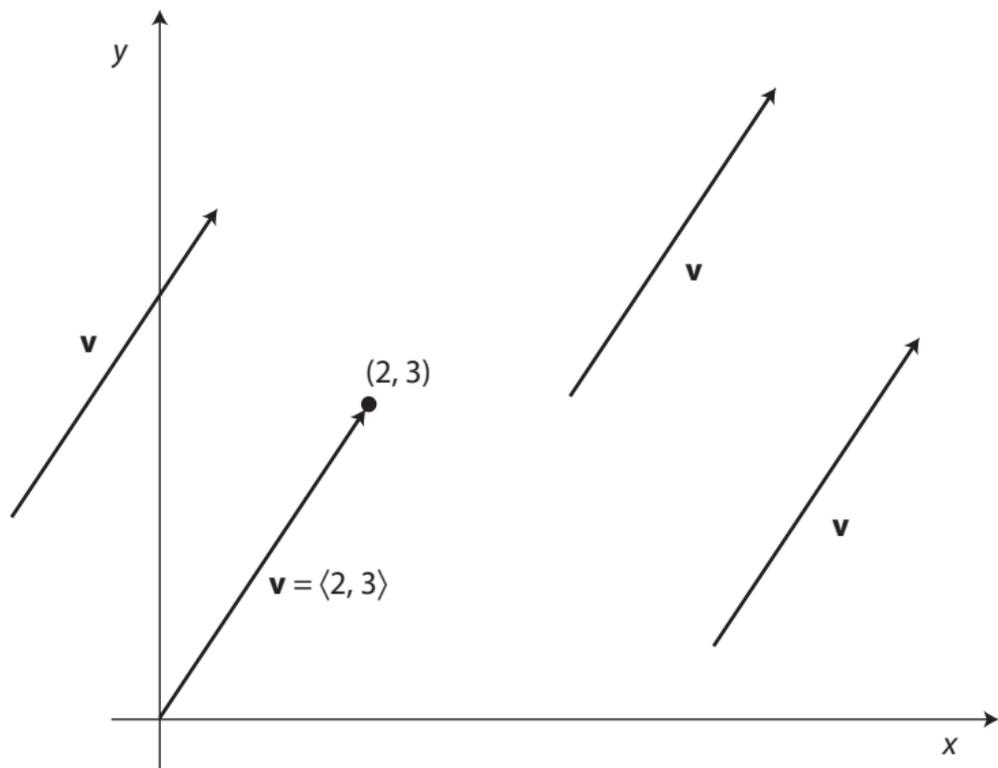


Figure 4.3: A vector has a direction and a length but it can start anywhere.

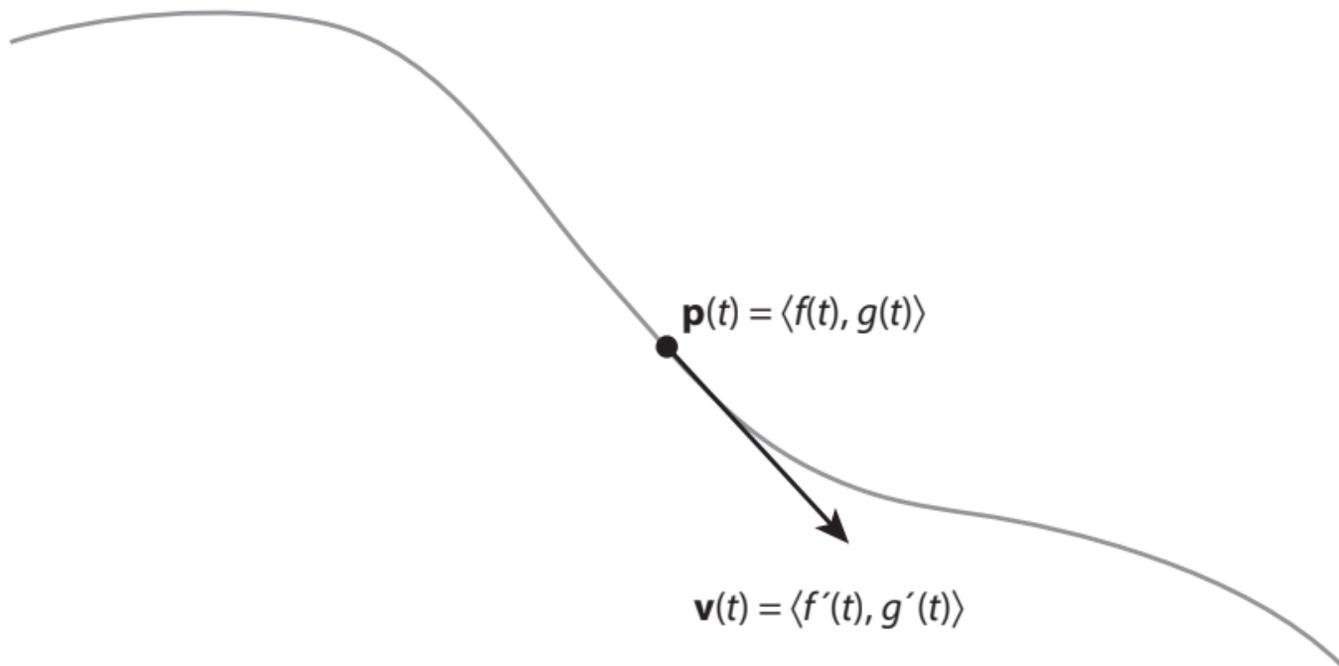


Figure 4.4: The tangent vector is given by $\mathbf{v}(t) = \langle f'(t), g'(t) \rangle$.

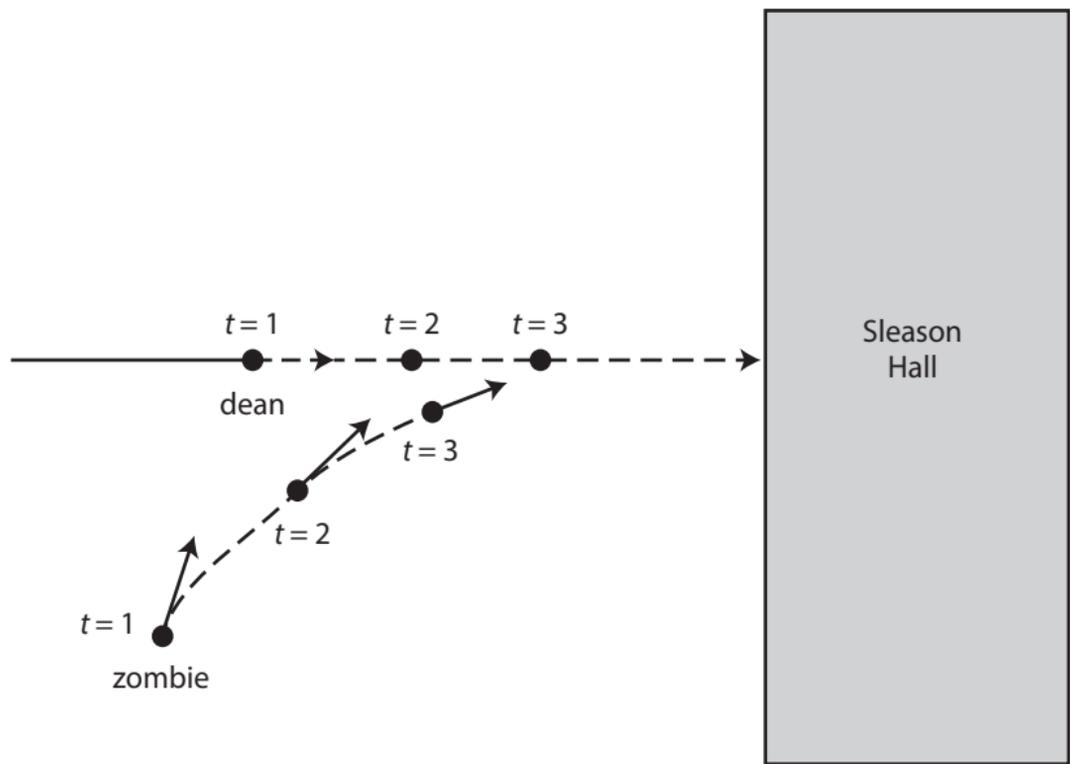


Figure 4.5: The zombie's tangent vector points toward the dean.

CHAPTER 4

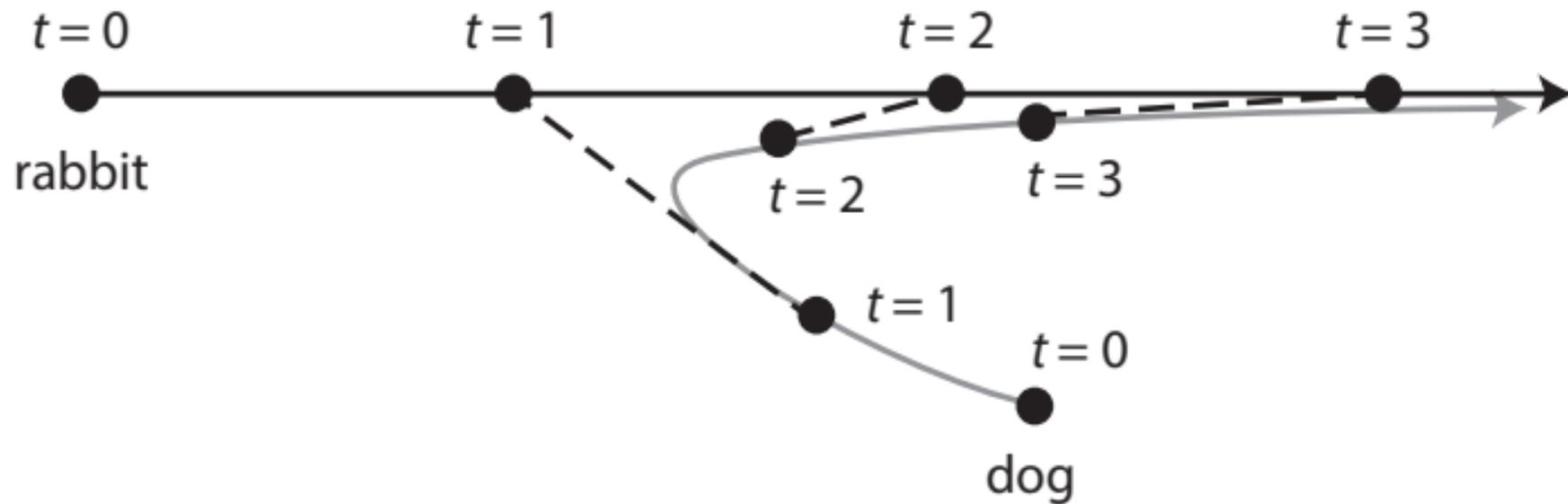


Figure 4.6: Slower dog chasing rabbit that passes nearby.

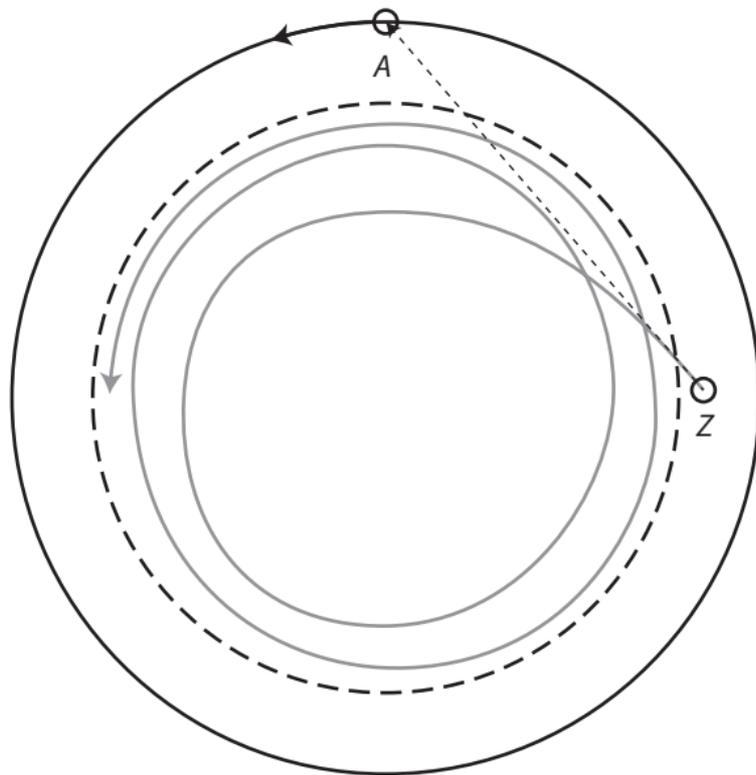


Figure 5.1: Zombie pursuing Angus, who rides his bike in a circle.

HOUR 8

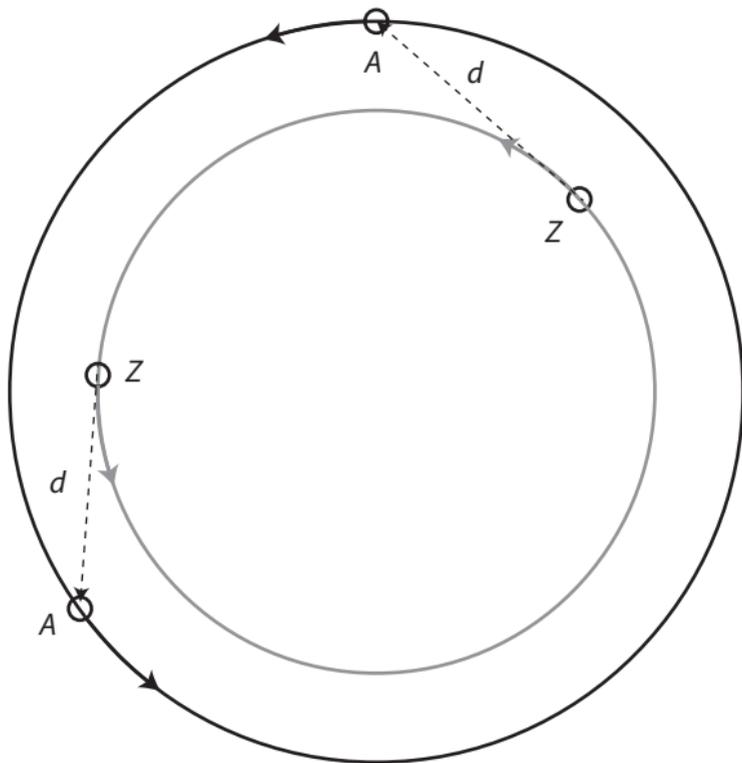


Figure 5.2: On the limit cycle, the zombie's distance to Angus never changes.

CHAPTER 6

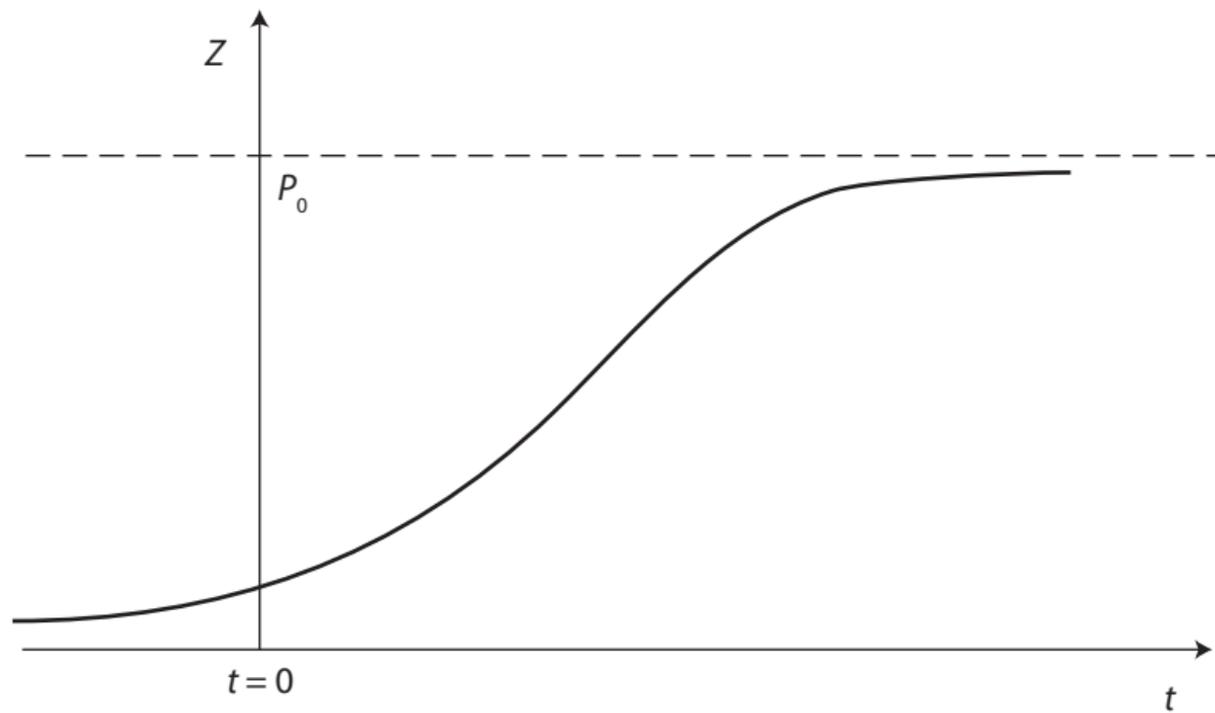


Figure 6.1: Logistic growth of the number of zombies.

HOUR 18

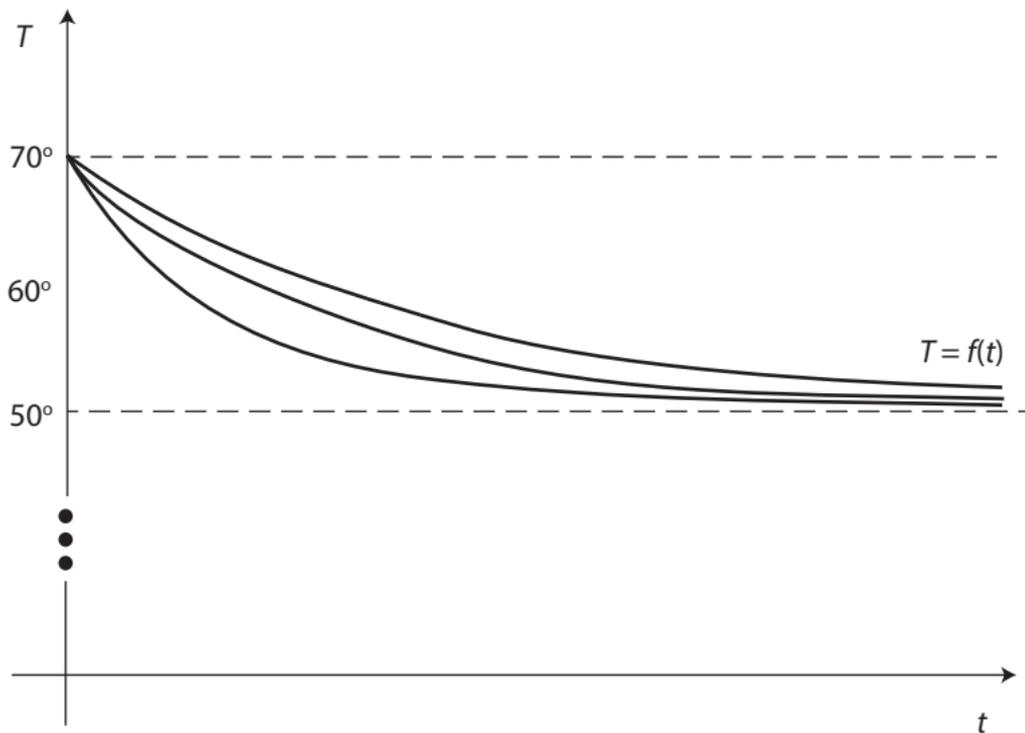


Figure 8.1: Zombie with initial temperature 70° cooling in a room of temperature 50° . Each curve corresponds to a different value of k .

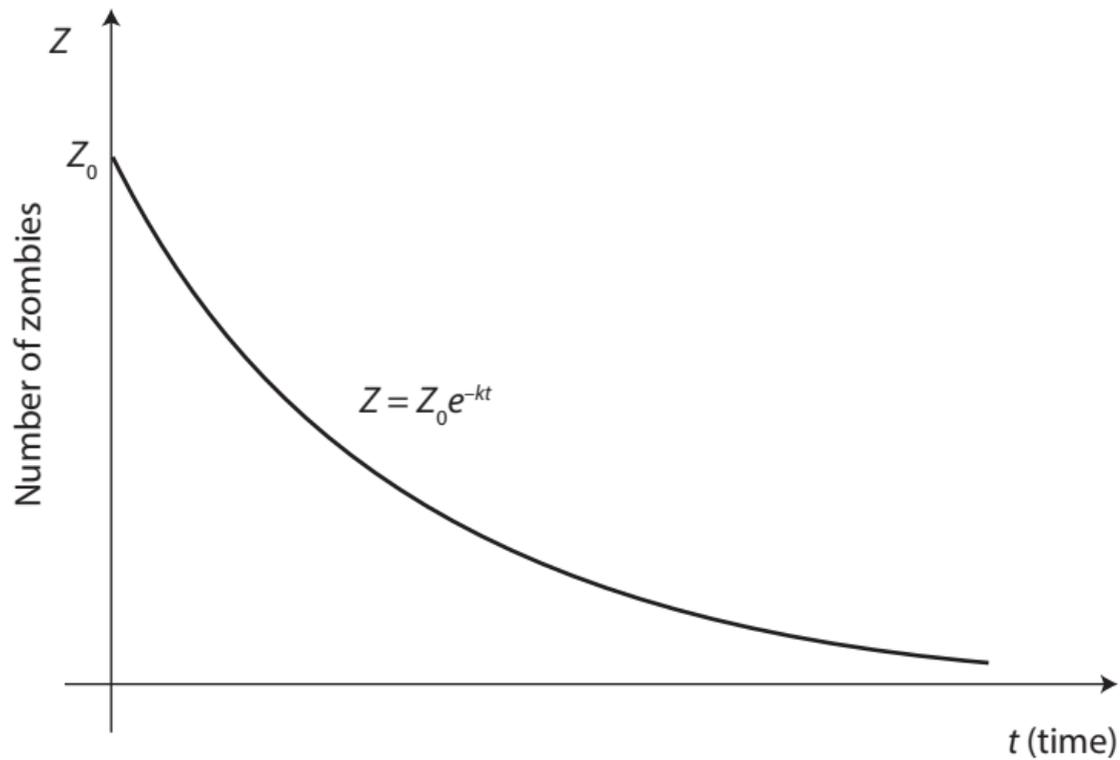


Figure 9.1: The decay in the number of zombies if there are no prey.

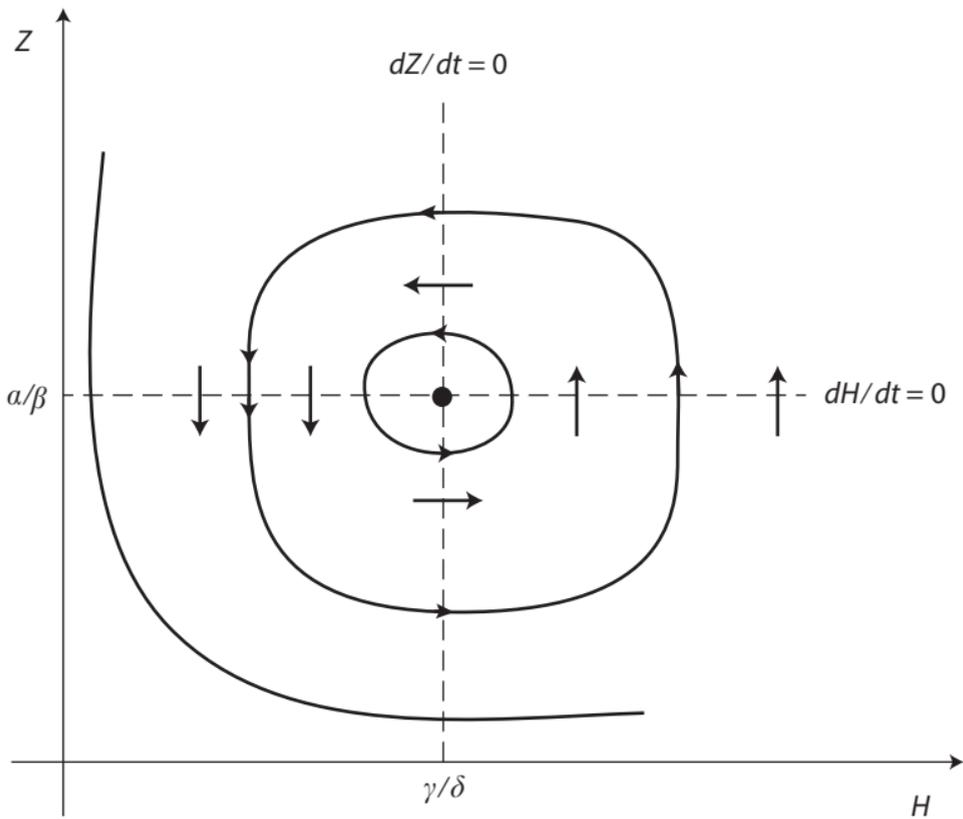


Figure 9.2: Curves in the HZ -plane.

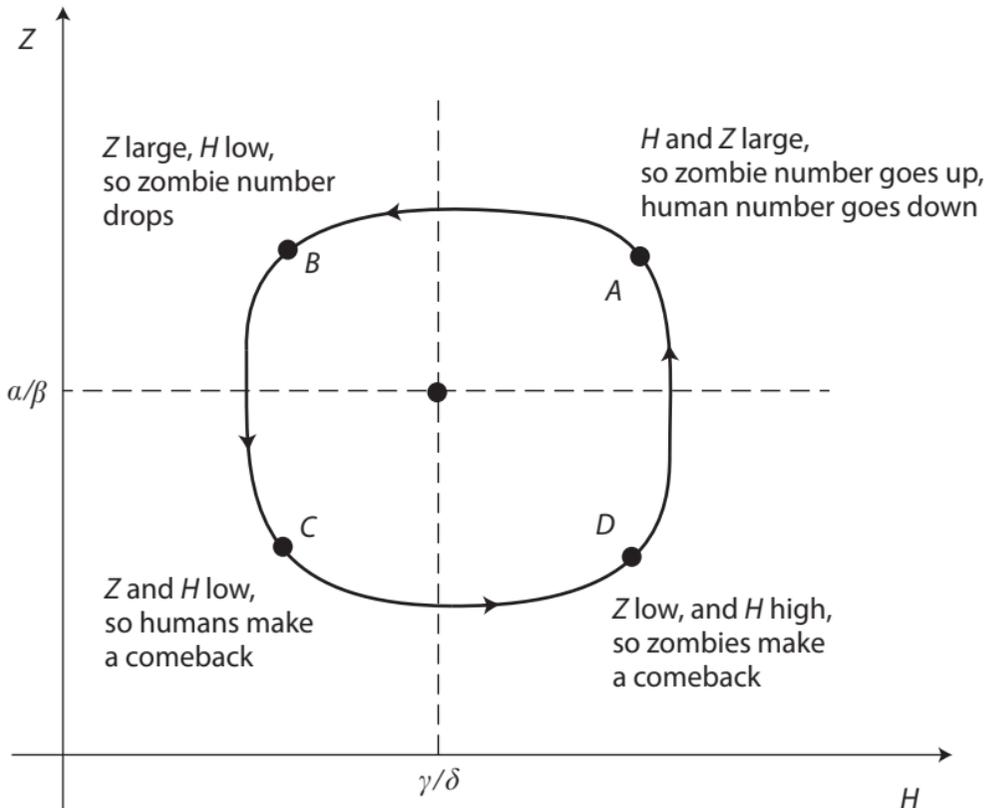


Figure 9.3: The predator-prey curve.

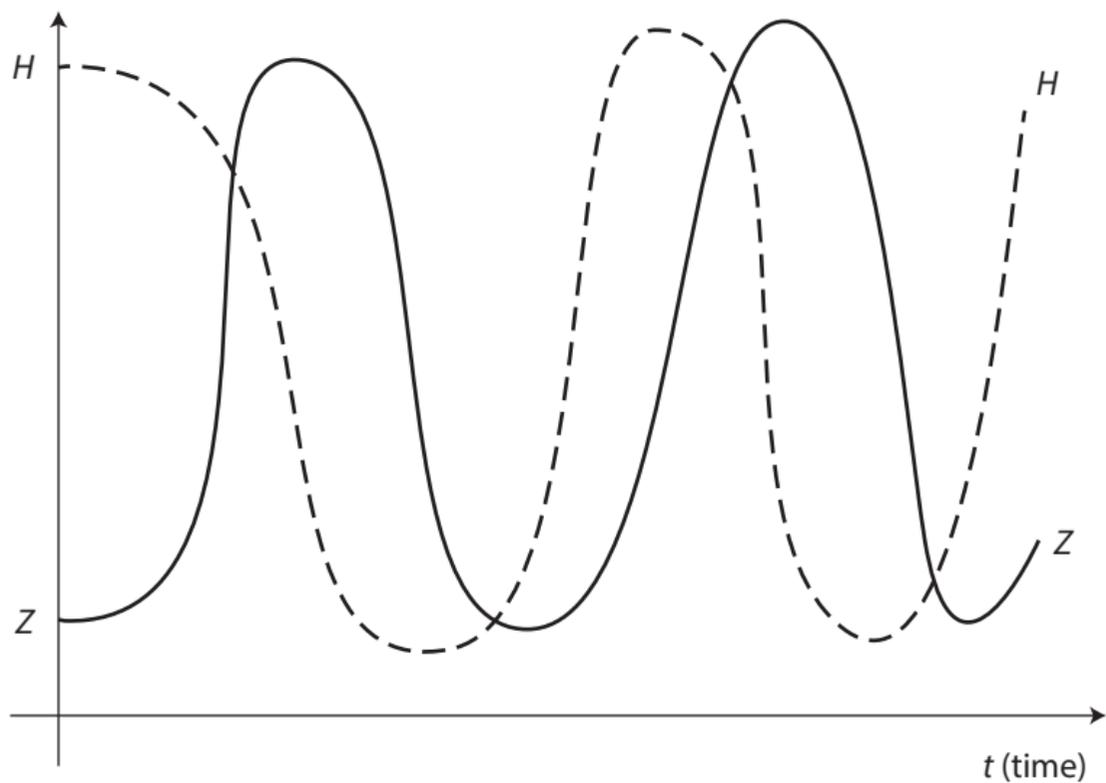


Figure 9.4: The oscillatory behavior of the human-zombie system.

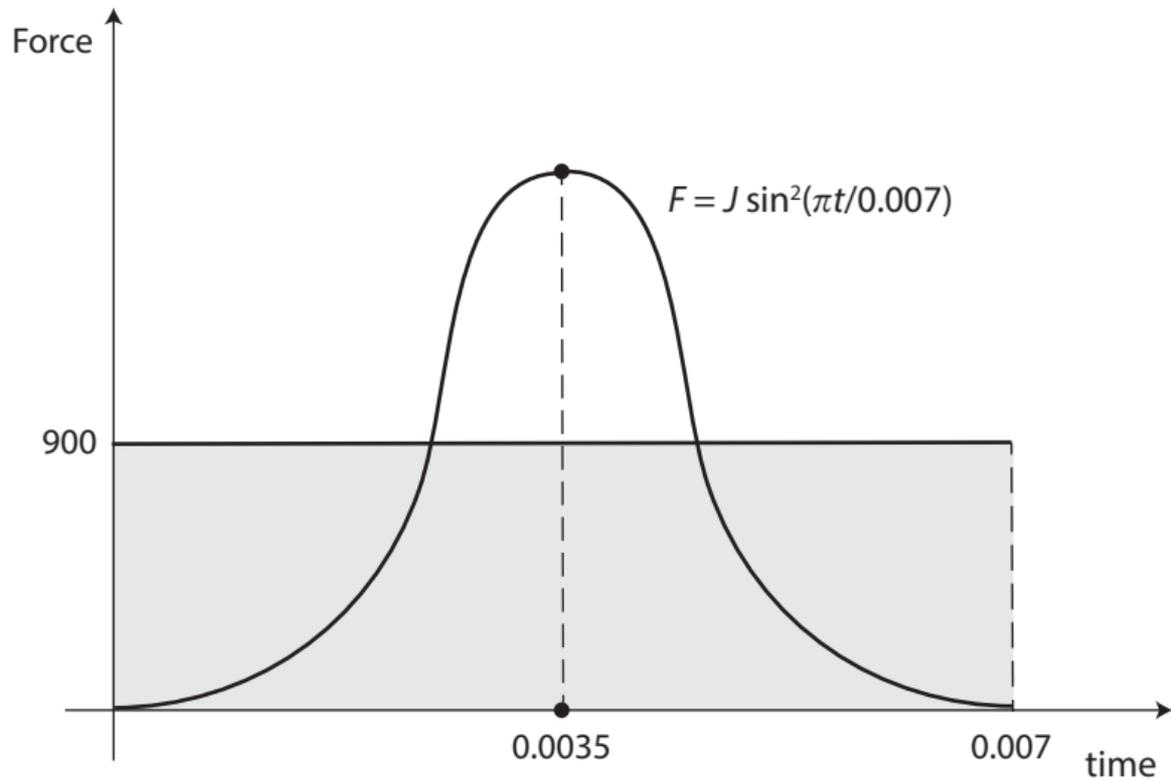


Figure A.1: The force curve.

So

$$\begin{aligned}\int J \sin^2 \left(\frac{\pi t}{0.007} \right) dt &= \int J \left(\frac{1 - \cos 2 \left(\frac{\pi t}{0.007} \right)}{2} \right) dt \\ &= J \left(\frac{1}{2} \int dt - \frac{1}{2} \int \cos \left(\frac{2\pi t}{0.007} \right) dt \right).\end{aligned}$$

The first integral is easy, but to do the second integral, we need to use u -substitution. We let $u = \frac{2\pi t}{0.007}$. Then $du = \frac{2\pi}{0.007} dt$, so $dt = \frac{0.007}{2\pi} du$ and we get:

$$\begin{aligned}\int \cos \left(\frac{2\pi t}{0.007} \right) dt &= \int \cos u \left(\frac{0.007}{2\pi} \right) du \\ &= \frac{0.007}{2\pi} \sin u + C \\ &= \frac{0.007}{2\pi} \sin \left(\frac{2\pi t}{0.007} \right) + C.\end{aligned}$$

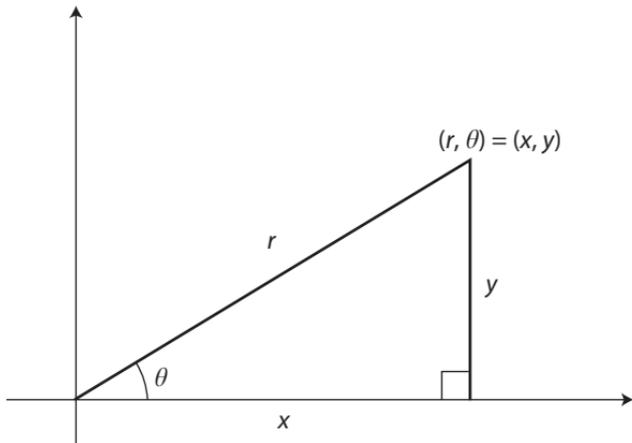
So we get

$$\int J \sin^2 \left(\frac{\pi t}{0.007} \right) dt = J \left(\frac{t}{2} - \frac{0.007}{4\pi} \sin \left(\frac{2\pi t}{0.007} \right) \right) + C.$$

Then, for the definite integral,

$$\begin{aligned}\int_0^{0.007} J \sin^2 \left(\frac{\pi t}{0.007} \right) dt &= J \left(\frac{t}{2} - \frac{0.007}{4\pi} \sin \left(\frac{2\pi t}{0.007} \right) \right) \Bigg|_0^{0.007} \\ &= J \left(\frac{0.007}{2} \right) = J(0.0035).\end{aligned}$$

CONTINUING THE CONVERSATIONS

**Figure A.2:** Polar coordinates.

Then

$$\begin{aligned} I^2 &= \left(\int_{-\infty}^{\infty} e^{-\frac{x^2}{2}} dx \right) \left(\int_{-\infty}^{\infty} e^{-\frac{y^2}{2}} dy \right) \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{x^2}{2}} e^{-\frac{y^2}{2}} dy dx \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{x^2+y^2}{2}} dy dx. \end{aligned}$$

APPENDIX A

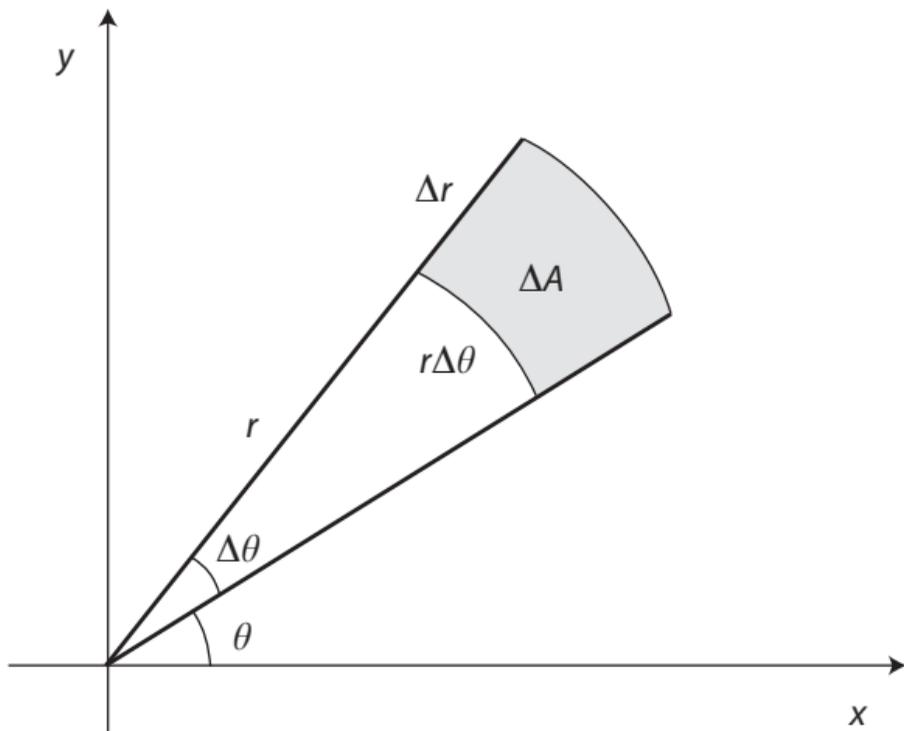


Figure A.3: Varying each coordinate.

“Yeah?”

“Yeah. Then, in the limit, instead of using $dA = dy dx$ in the integral, we replace it with $r dr d\theta$. So we get

$$\begin{aligned} I^2 &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{\frac{-x^2-y^2}{2}} dy dx \\ &= \int_0^{2\pi} \int_0^{\infty} e^{\frac{-r^2}{2}} r dr d\theta. \end{aligned}$$

Now, we can do the inner r -integral by u -substitution, with $u = \frac{-r^2}{2}$. Then $du = -r dr$:

$$\begin{aligned} I^2 &= \int_0^{2\pi} \int_0^{\infty} e^{\frac{-r^2}{2}} r dr d\theta \\ &= \int_0^{2\pi} \int_0^{-\infty} e^u (-1) du d\theta \\ &= \int_0^{2\pi} (-e^u) \Big|_0^{-\infty} d\theta \\ &= \int_0^{2\pi} (0 - (-1)) d\theta \\ &= 2\pi. \end{aligned}$$

So $I^2 = 2\pi$ and since I is clearly positive, $I = \sqrt{2\pi}$. That’s what we wanted to show.”

“That’s a pretty cool trick,” said Angus.

CONTINUING THE CONVERSATIONS

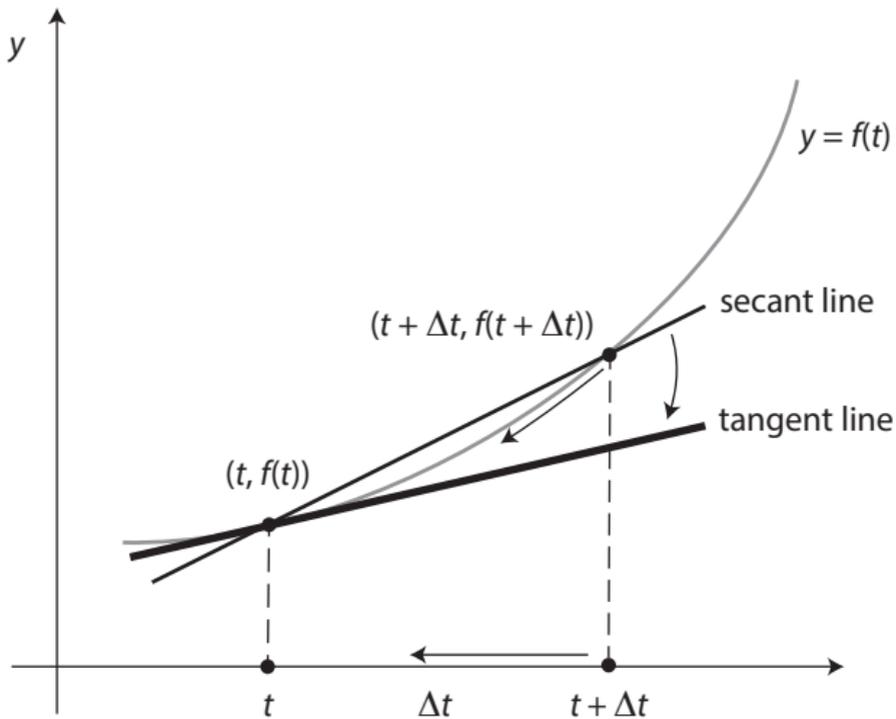


Figure A.4: The definition of the derivative.

APPENDIX A

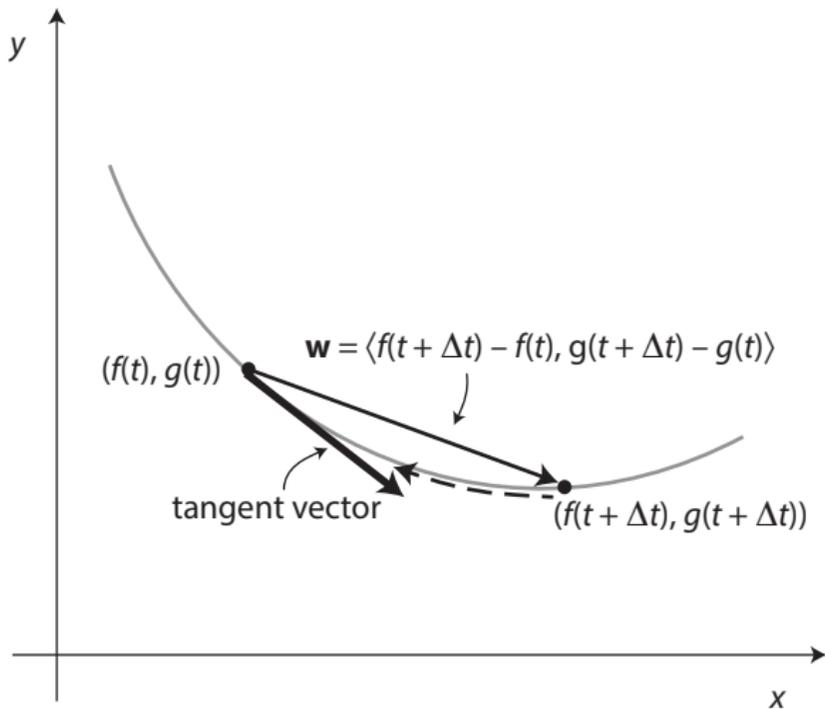


Figure A.5: The velocity vector is tangent.

APPENDIX A

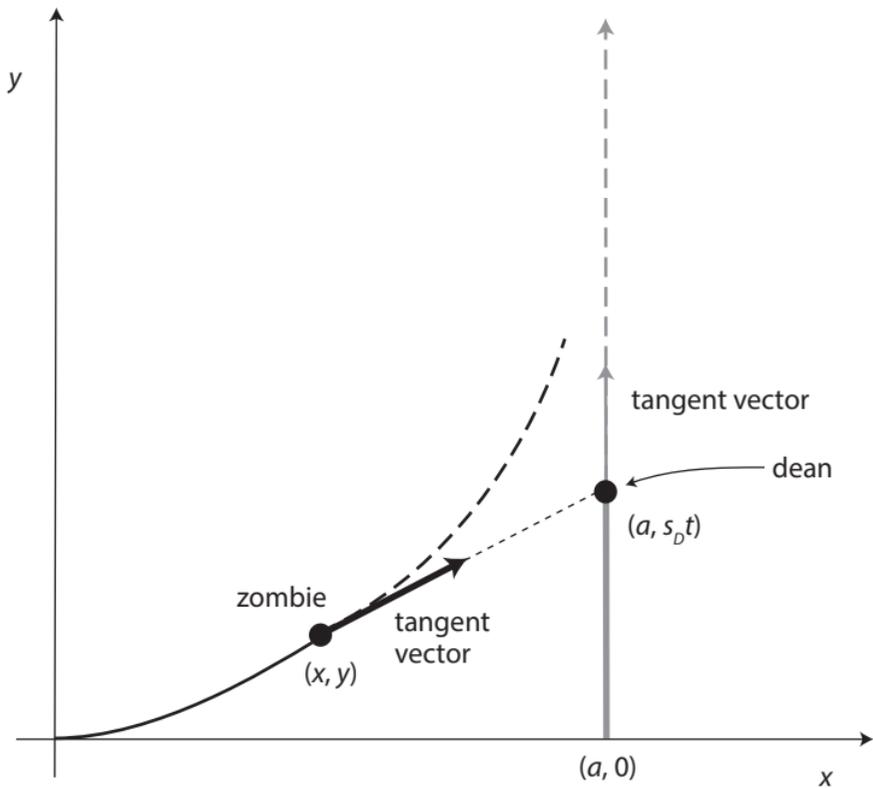


Figure A.6: Zombie chasing the dean.

CONTINUING THE CONVERSATIONS

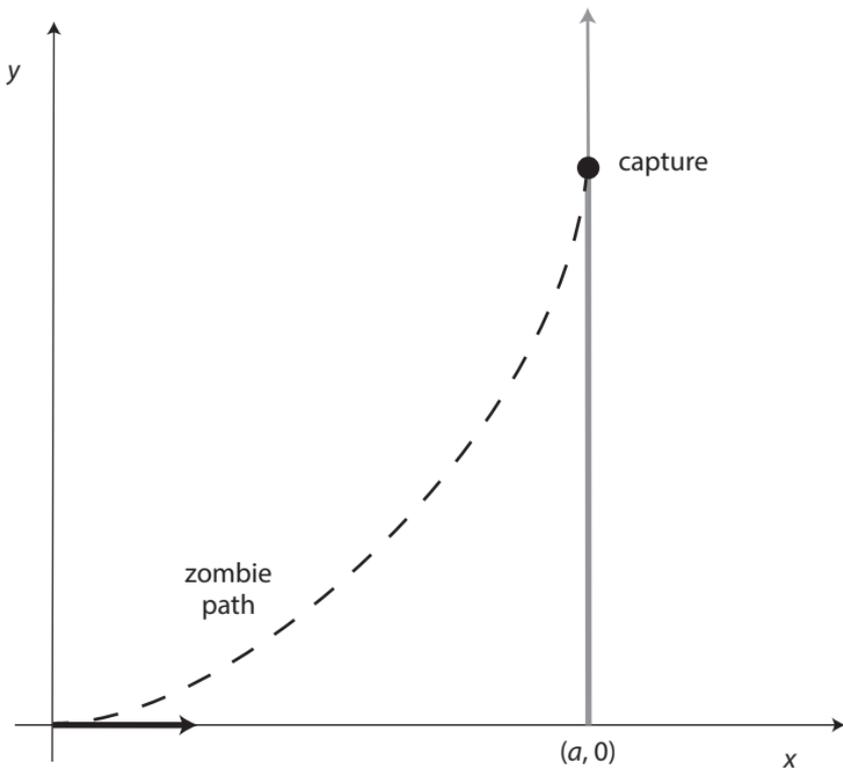


Figure A.7: Zombie catches the dean from directly behind.

CONTINUING THE CONVERSATIONS

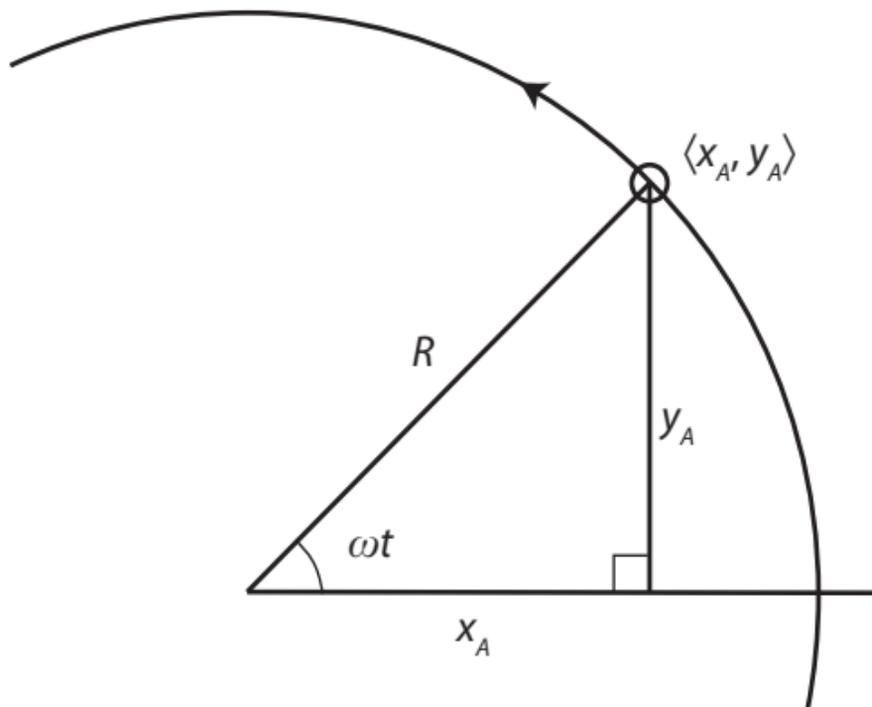


Figure A.8: Angus's path in circular pursuit.

APPENDIX A

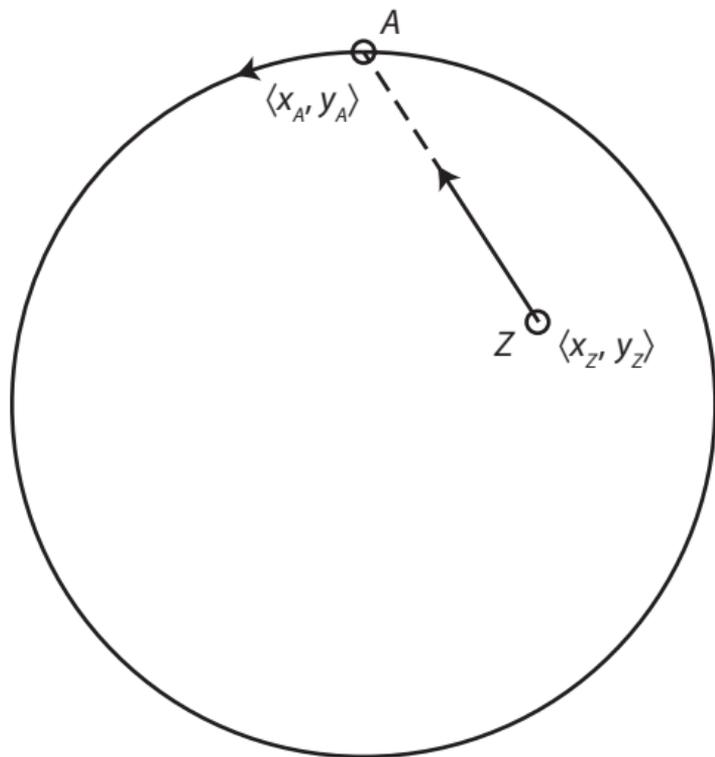


Figure A.9: Determining the zombie's path in circular pursuit.

“By considering each component separately, we obtain two coupled differential equations:

$$\frac{dx_Z}{dt} = s_Z \frac{x_A - x_Z}{\sqrt{(x_A - x_Z)^2 + (y_A - y_Z)^2}},$$

$$\frac{dy_Z}{dt} = s_Z \frac{y_A - y_Z}{\sqrt{(x_A - x_Z)^2 + (y_A - y_Z)^2}}.$$

“Adding in the expressions for x_A and y_A , we obtain:

$$\frac{dx_Z}{dt} = s_Z \frac{R \cos(\omega t) - x_Z}{\sqrt{(R \cos(\omega t) - x_Z)^2 + (R \sin(\omega t) - y_Z)^2}},$$

$$\frac{dy_Z}{dt} = s_Z \frac{R \sin(\omega t) - y_Z}{\sqrt{R \cos(\omega t) - x_Z)^2 + (R \sin(\omega t) - y_Z)^2}}.$$

“So these are the differential equations that need to be satisfied by the zombie’s path. But they’re too complicated to get an analytic solution. Instead, you get a computer to draw the path.”



VII. Logistic growth (continued from p. 84)

“As Jessie explained,” I said, “a disease often spreads so that the rate of spread is proportional to the number infected and the number not yet infected. So we have a differential equation that looks like this:

$$\frac{dZ}{dt} = kZ(P_0 - Z)$$

where P_0 is the initial population.”

Since Z can be anything, we can set it to whatever we want. So we set it to 0, and we find $1 = AP_0$. So $A = 1/P_0$. Plugging that in, the equation becomes

$$1 = \frac{1}{P_0}(P_0 - Z) + BZ,$$

$$1 = 1 + Z \left(\frac{-1}{P_0} + B \right),$$

$$\frac{-1}{P_0} + B = 0,$$

$$B = \frac{1}{P_0}."$$

"And so that means what?"

"That means we now know

$$\frac{1}{Z(P_0 - Z)} = \frac{1}{P_0Z} + \frac{1}{P_0(P_0 - Z)}.$$

So, now we can integrate:

$$\int \frac{dZ}{Z(P_0 - Z)} = \int \frac{dZ}{P_0Z} + \int \frac{dZ}{P_0(P_0 - Z)}.$$

"We can do both these integrals. We get

$$\begin{aligned} \int \frac{dZ}{Z(P_0 - Z)} &= \frac{1}{P_0} \ln Z - \frac{1}{P_0} \ln(P_0 - Z) + C \\ &= \frac{1}{P_0} \ln \frac{Z}{P_0 - Z} + C. \end{aligned}$$

“Then $\int \frac{dZ}{Z(P_0-Z)} = \int k dt$ becomes

$$\frac{1}{P_0} \ln \frac{Z}{P_0 - Z} = kt + C,$$

$$\ln \frac{Z}{P_0 - Z} = P_0 kt + P_0 C,$$

$$\frac{Z}{P_0 - Z} = e^{P_0 kt + P_0 C}.$$

Notice $e^{P_0 C}$ is just a constant. Let's call it J . Then we have

$$\frac{Z}{P_0 - Z} = J e^{P_0 kt}.$$

Now solve for Z .”

“How do we do that?” asked Angus.

“We multiply through by $P_0 - Z$.”

$$Z = (P_0 - Z) J e^{P_0 kt}$$

$$Z = P_0 J e^{P_0 kt} - Z J e^{P_0 kt}$$

$$Z(1 + J e^{P_0 kt}) = P_0 J e^{P_0 kt}$$

$$Z = \frac{P_0 J e^{P_0 kt}}{1 + J e^{P_0 kt}}$$

“So that is our function that gives Z as a function of time t .”

“In this case, we’ll set

$$u = t \quad \text{and} \quad dv = e^{-\alpha t} dt.$$

$$\text{Then } du = \frac{du}{dt} dt = dt,$$

$$\text{and } v = \int dv = \int e^{-\alpha t} dt = \frac{-1}{\alpha} e^{-\alpha t}.$$

So according to the formula,

$$\tau = \alpha \int_0^{\infty} t e^{-\alpha t} dt = \alpha \left(t \frac{-1}{\alpha} e^{-\alpha t} \right]_0^{\infty} - \int_0^{\infty} \frac{-1}{\alpha} e^{-\alpha t} dt \Big).”$$

“But you still have an integral to do,” said Jessie.

“Yes, but it’s easier than the first one. We can do it directly, and we get

$$\begin{aligned} \tau &= \alpha \left(t \frac{-1}{\alpha} e^{-\alpha t} - \frac{1}{\alpha^2} e^{-\alpha t} \right) \Big]_0^{\infty} \\ &= -\frac{t}{e^{\alpha t}} - \frac{1}{\alpha e^{\alpha t}} \Big]_0^{\infty} \\ &= \lim_{b \rightarrow \infty} -\frac{t}{e^{\alpha t}} - \frac{1}{\alpha e^{\alpha t}} \Big]_0^b \\ &= \lim_{b \rightarrow \infty} -\frac{b}{e^{\alpha b}} - \frac{1}{\alpha e^{\alpha b}} - \left(-\frac{1}{\alpha} \right).” \end{aligned}$$

“Wait,” said Jessie. “How do we take the limit of $-\frac{b}{e^{\alpha b}}$? It gives us $\frac{\infty}{\infty}$.”

I smiled. “We use L’Hôpital’s Rule.”

“Why are you smiling?”

APPENDIX A

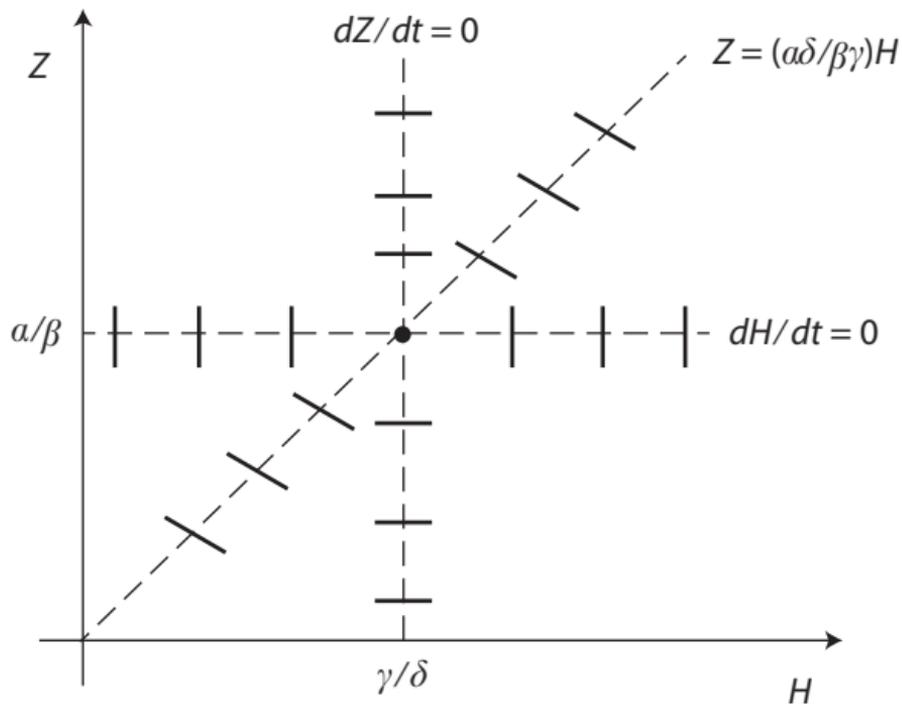


Figure A.10: The slope field in the HZ -plane.

CONTINUING THE CONVERSATIONS

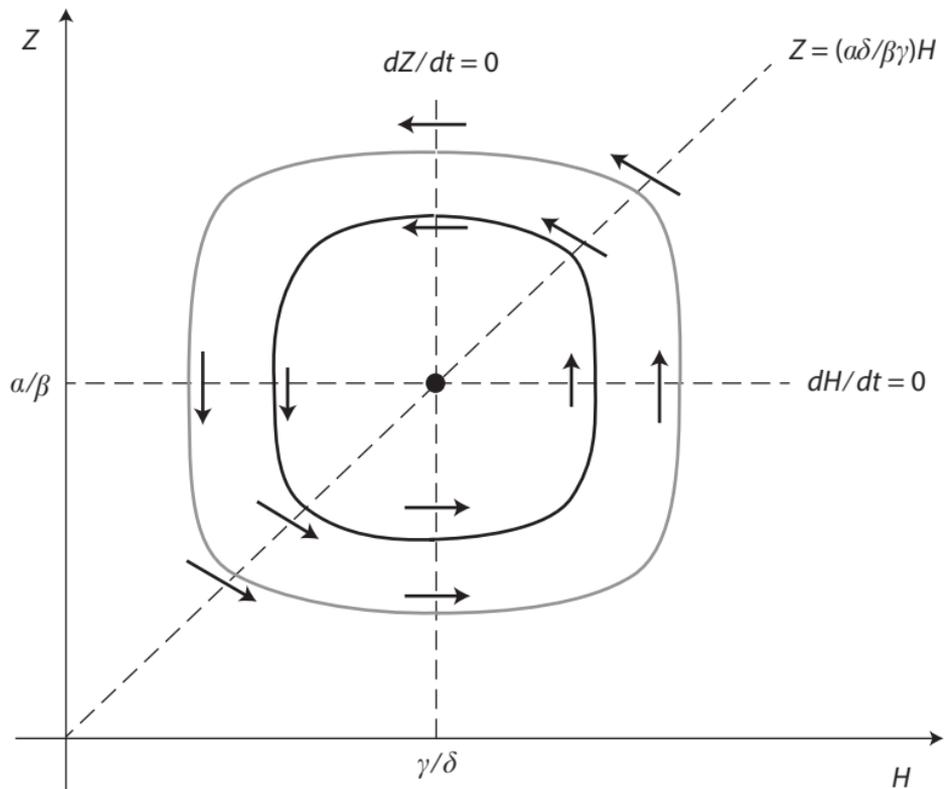


Figure A.11: Curves in the slope field in the HZ -plane.

APPENDIX B

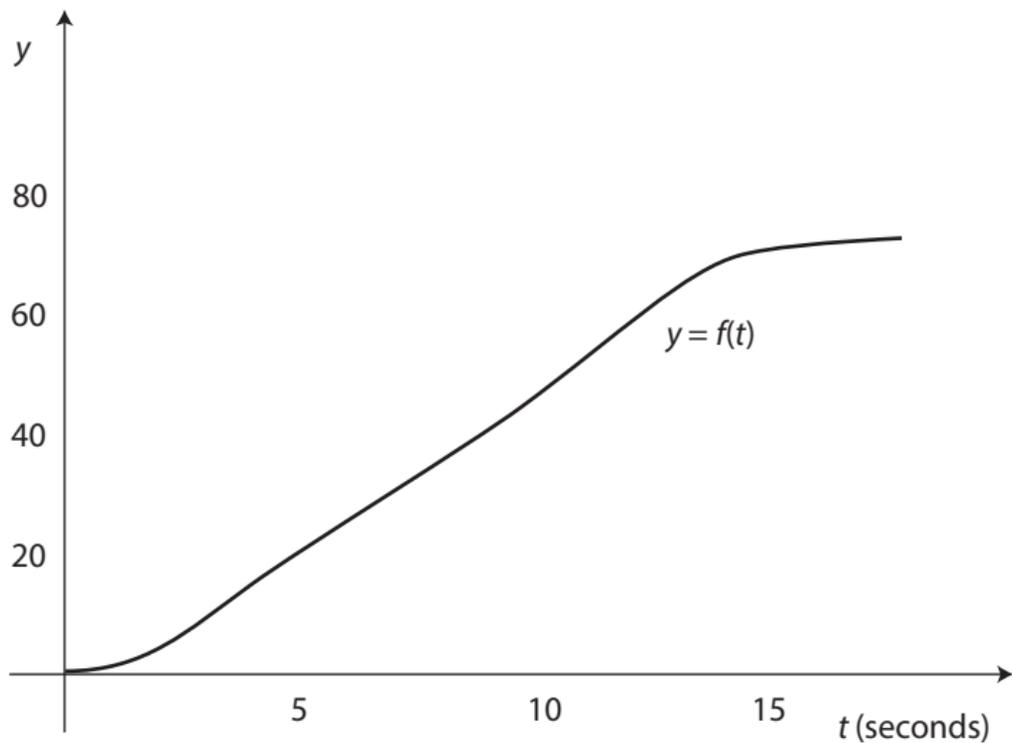


Figure B.1: Distance travelled while running from a zombie as a function of time.

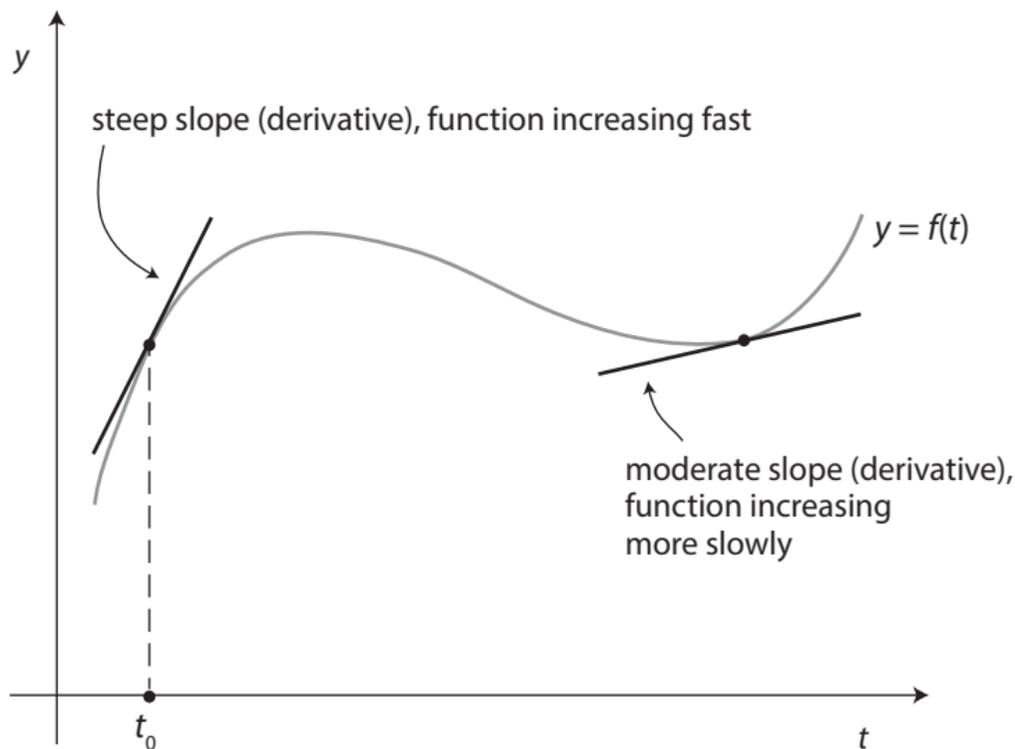


Figure B.2: The slope of the tangent line to a graph measures how steep the graph is at that point.

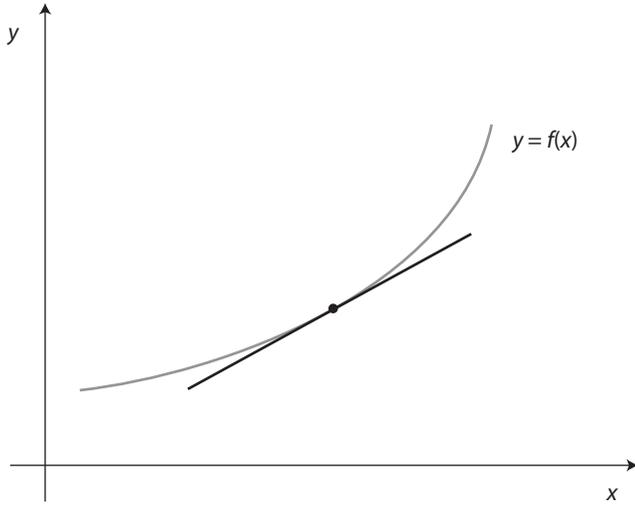


Figure B.3: Positive derivative means positive slope, which means the function is increasing.

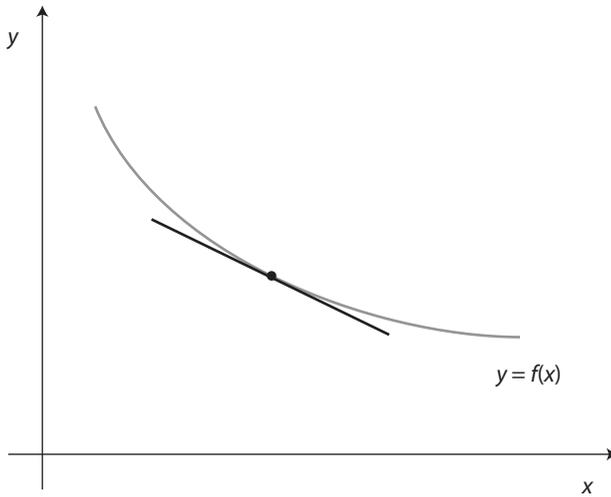


Figure B.4: Negative derivative means negative slope, which means the function is decreasing.

APPENDIX B

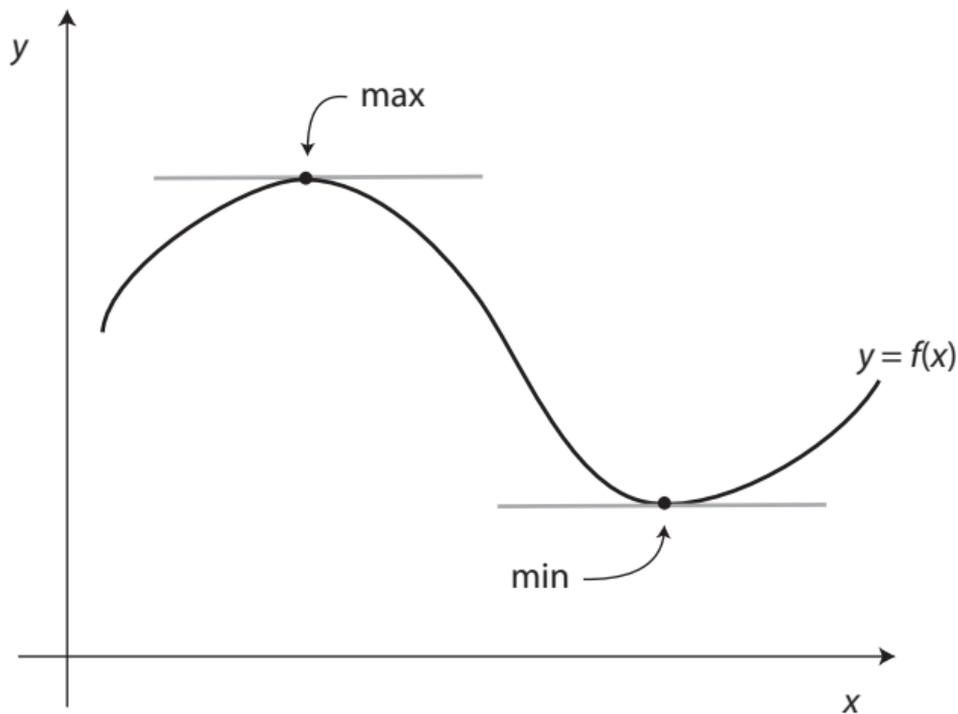


Figure B.5: At a local maximum or minimum, the derivative is 0.

APPENDIX B

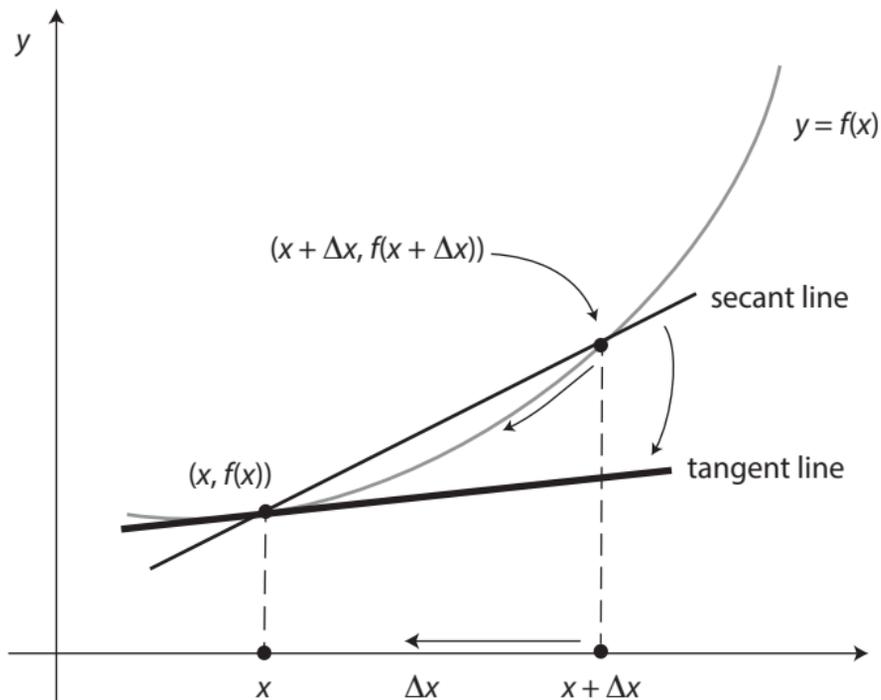


Figure B.6: The slope of the secant line approximates the slope of the tangent line.

APPENDIX B

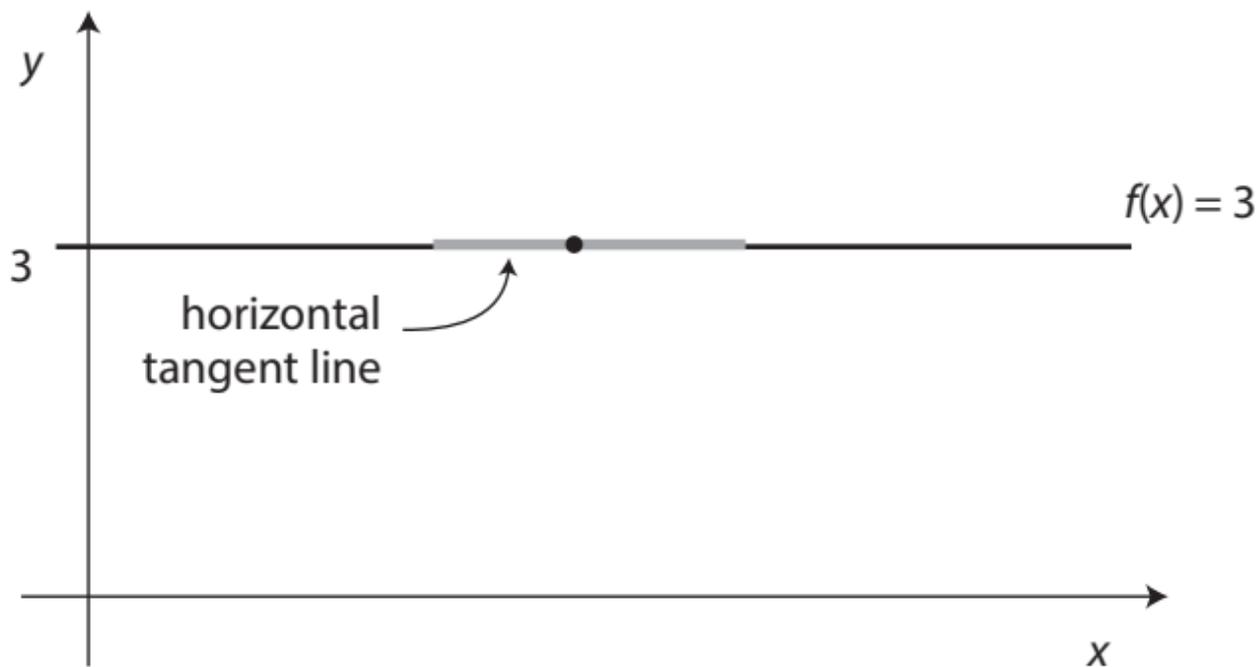


Figure B.7: For a constant function, slope and hence derivative is 0 everywhere.

APPENDIX B

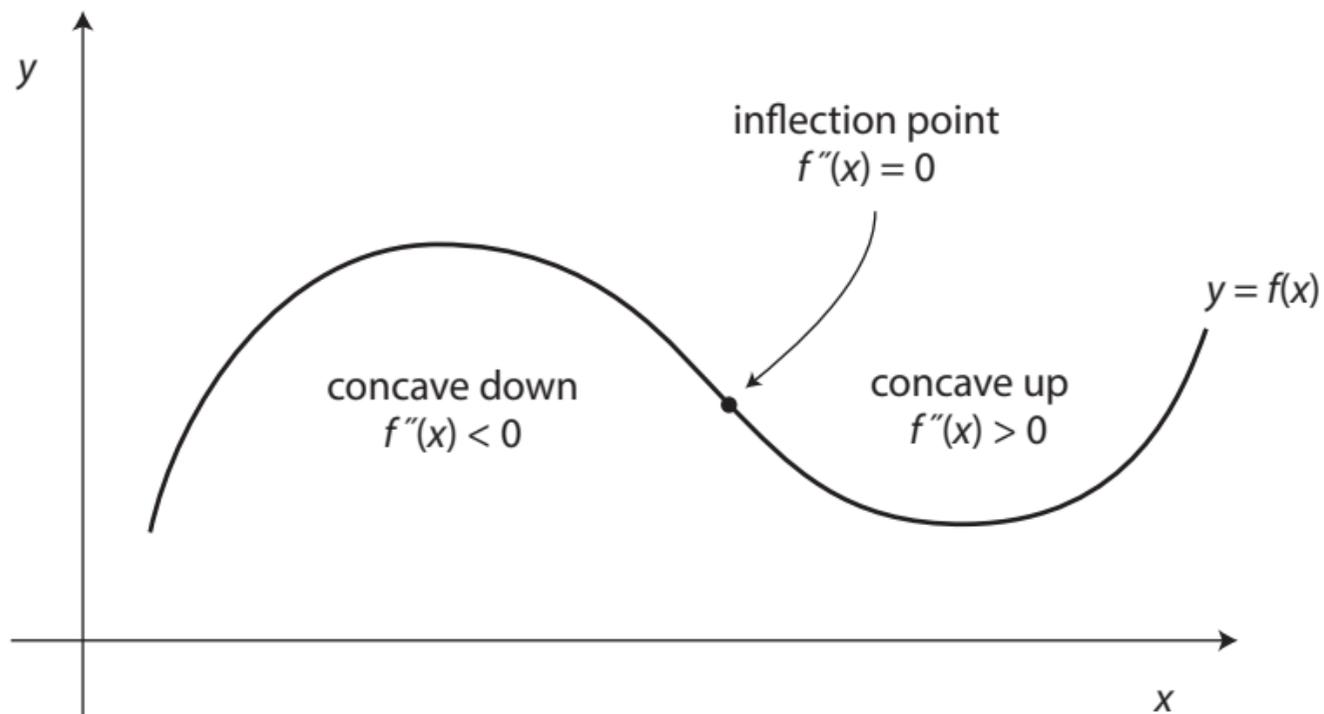


Figure B.8: $f''(x)$ determines concavity.

A BRIEF REVIEW OF CALCULUS

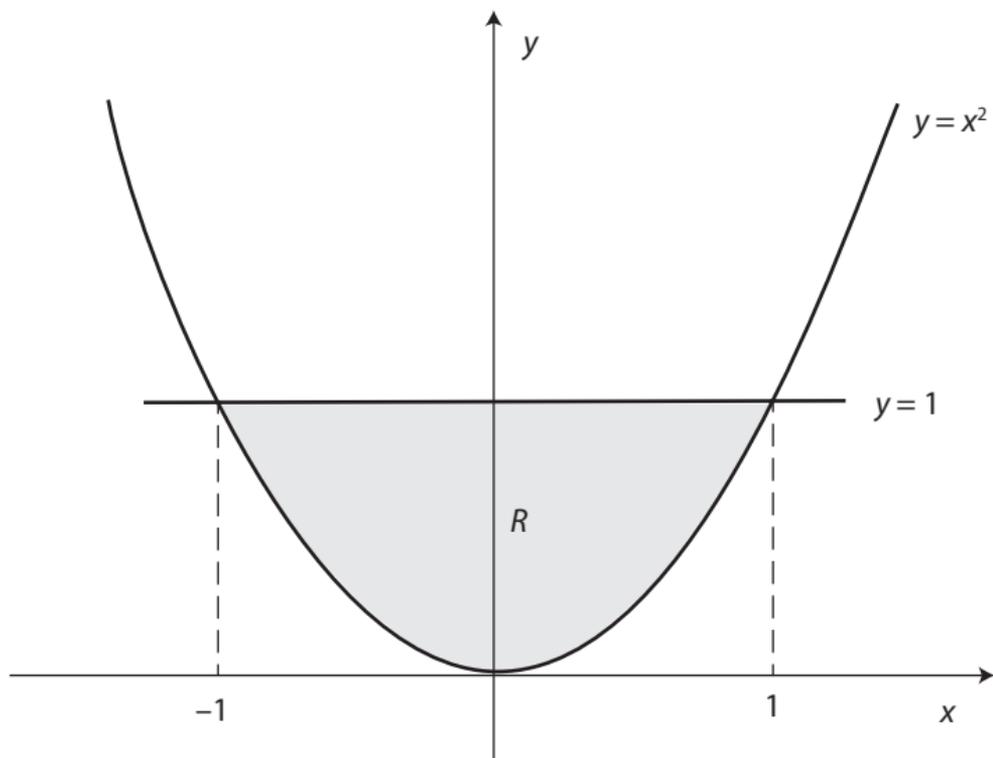


Figure B.9: Integrating over this region R .

APPENDIX B

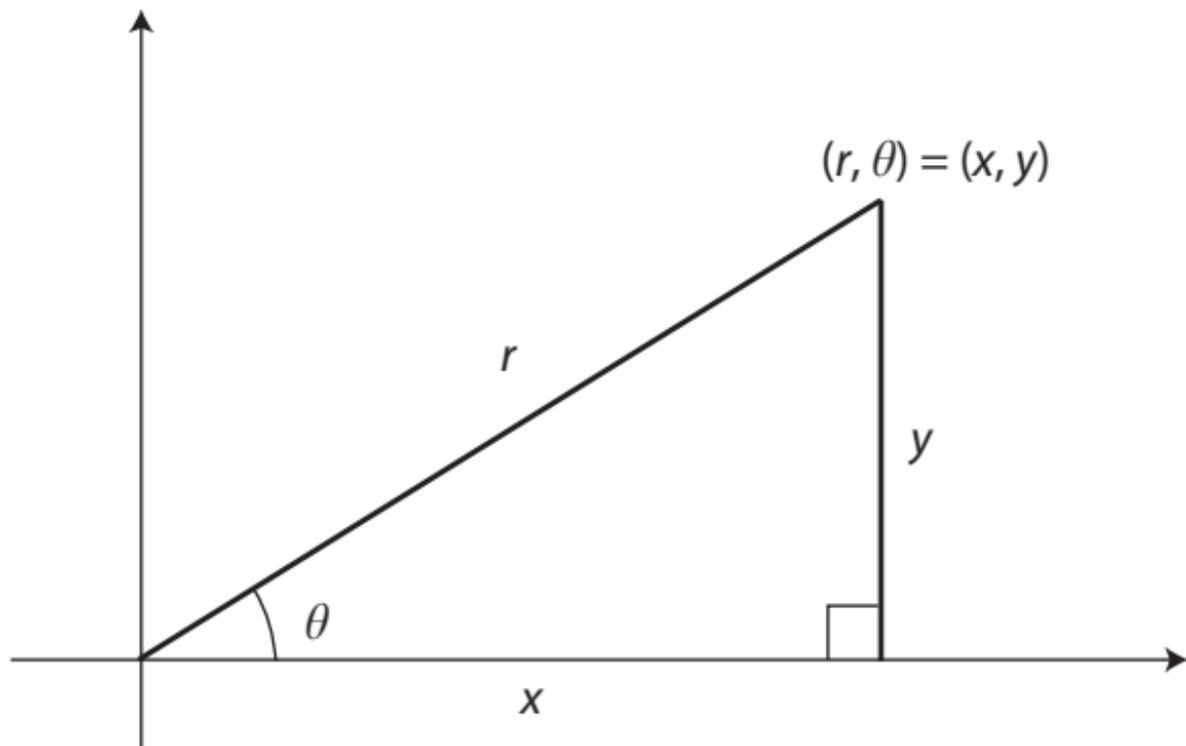


Figure B.10: Polar coordinates.

A BRIEF REVIEW OF CALCULUS

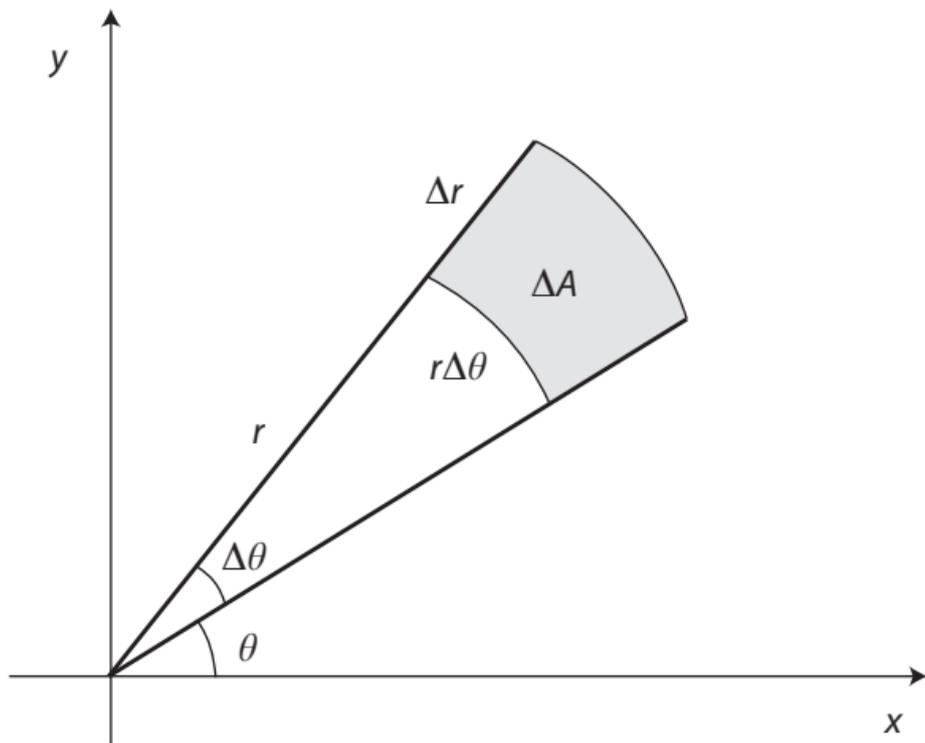


Figure B.11: A small area in polar coordinates.

A BRIEF REVIEW OF CALCULUS

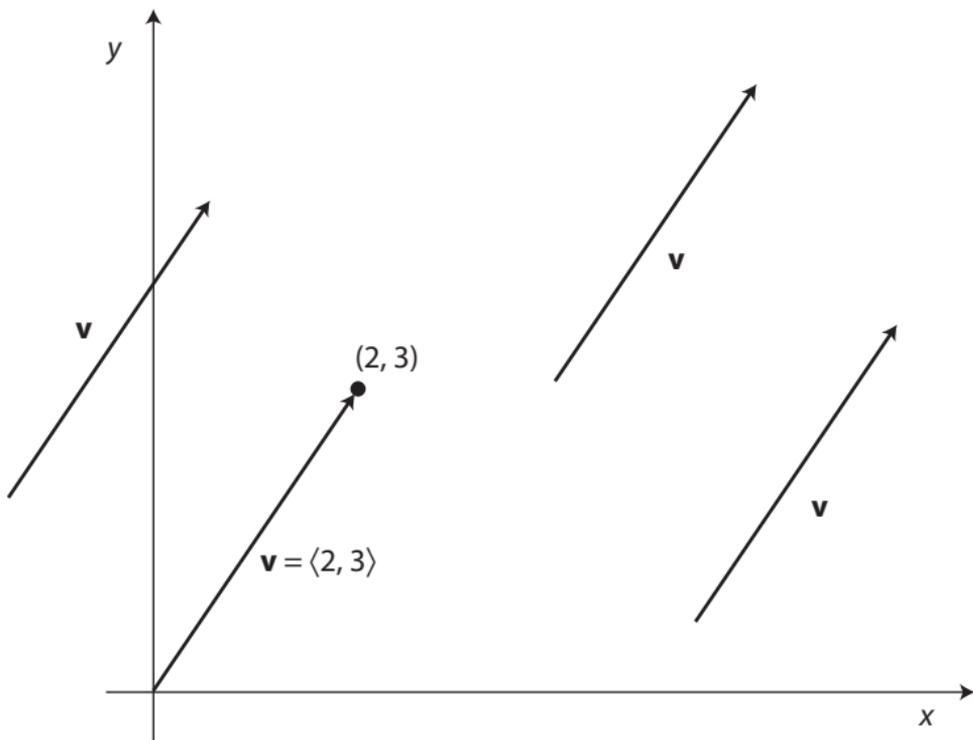


Figure B.12: The vector $\mathbf{v} = \langle 3, 2 \rangle$. It has direction and length.

A BRIEF REVIEW OF CALCULUS

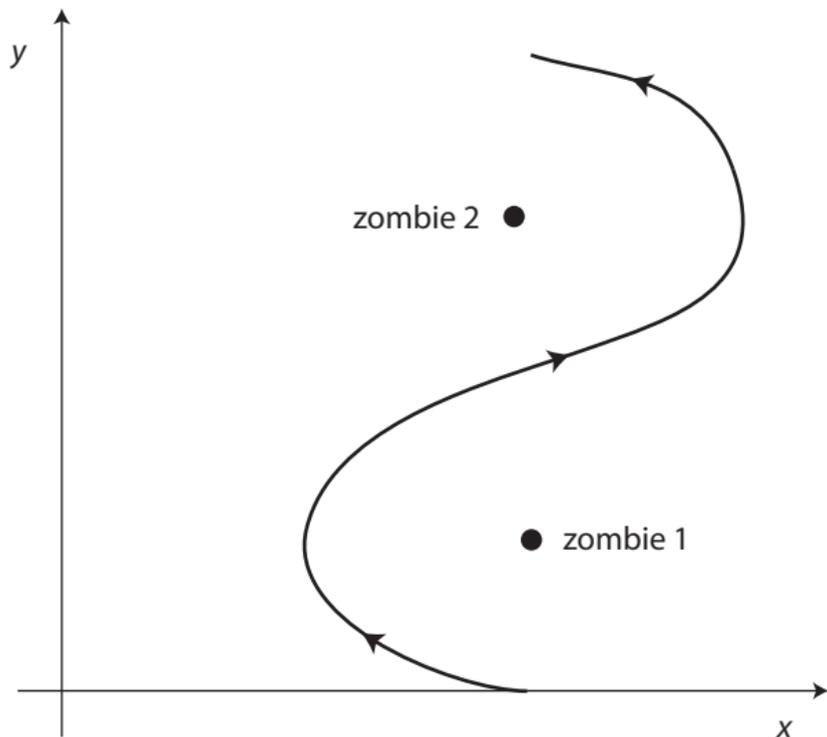


Figure B.13: The path taken by the provost to avoid zombies.

A BRIEF REVIEW OF CALCULUS

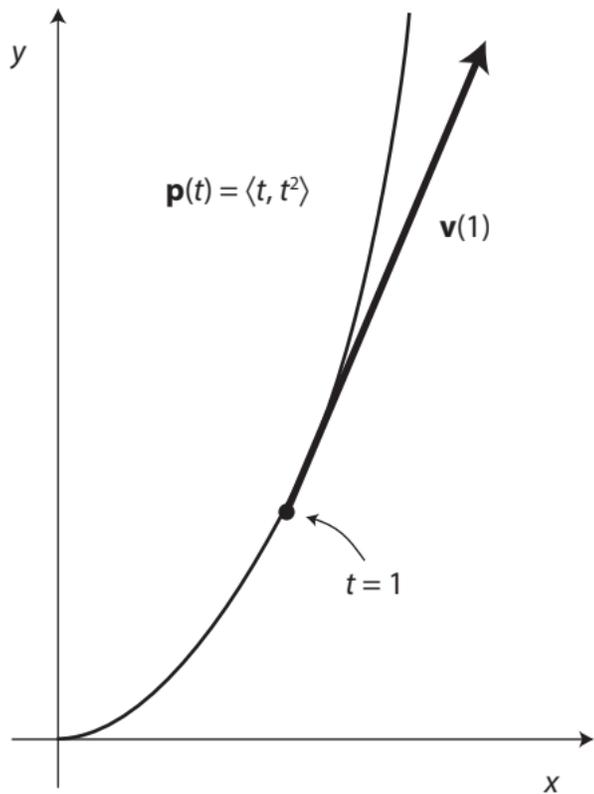


Figure B.14: The provost's parabolic path showing a velocity vector.